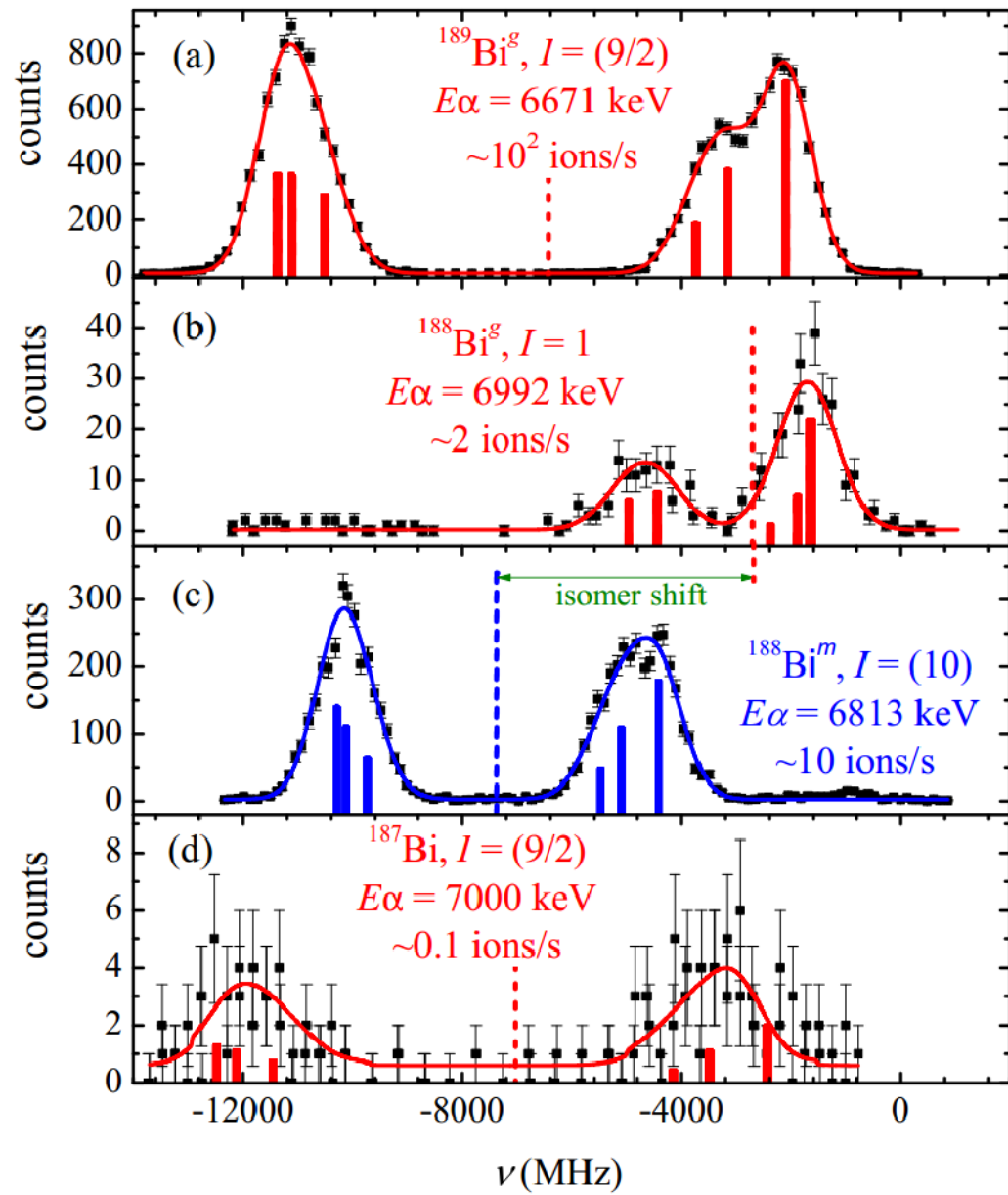
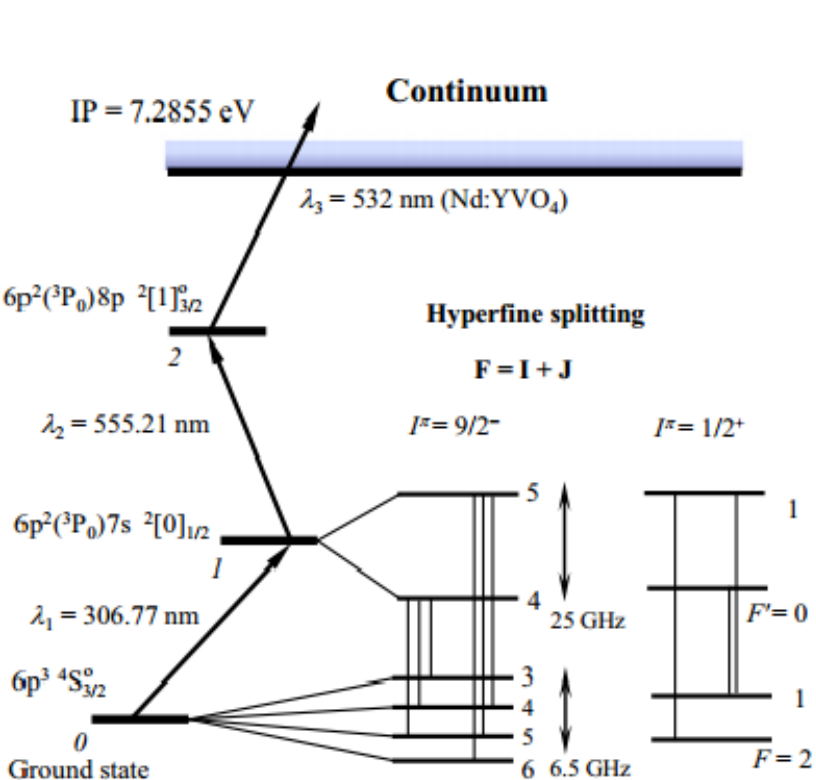


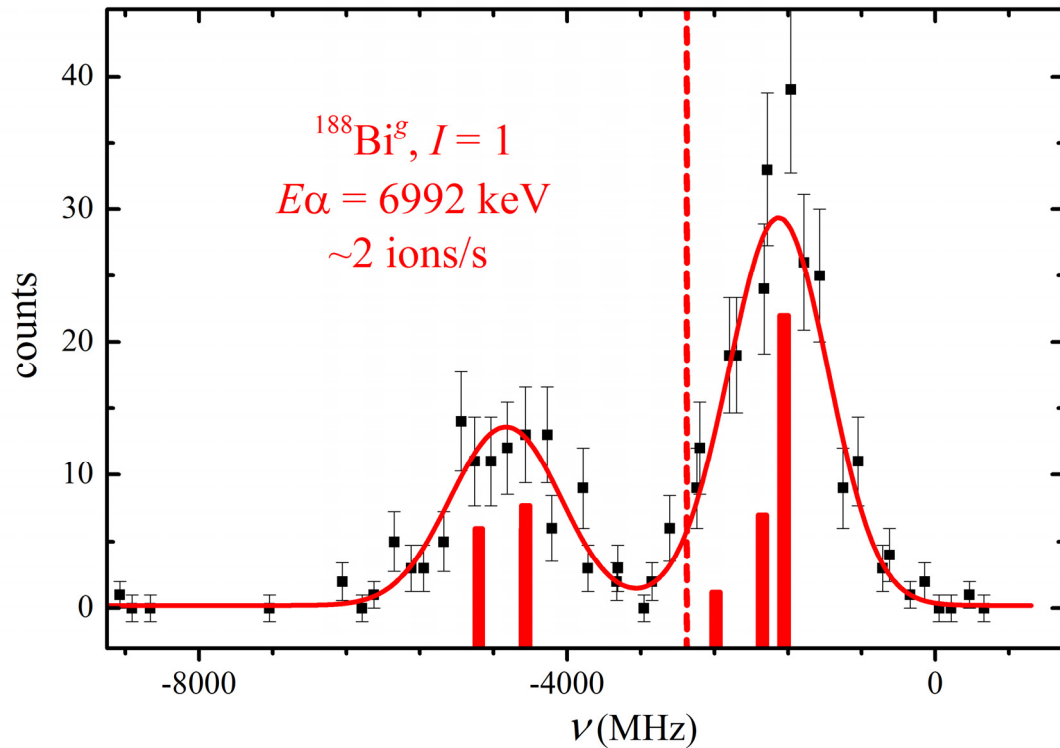
**Эксперименты с лазерным ионным
источником (ISOLDE, CERN;
Windmill-ISOLTRAP-RILIS-IDS collaboration)**

А. Е. Барзах, П. Л. Молканов, В. Н. Пантелеев,
М. Д. Селиверстов, Д. В. Федоров

IS and hfs for neutron-deficient Bi isotopes



Spin of $^{188}\text{Bi}^g$



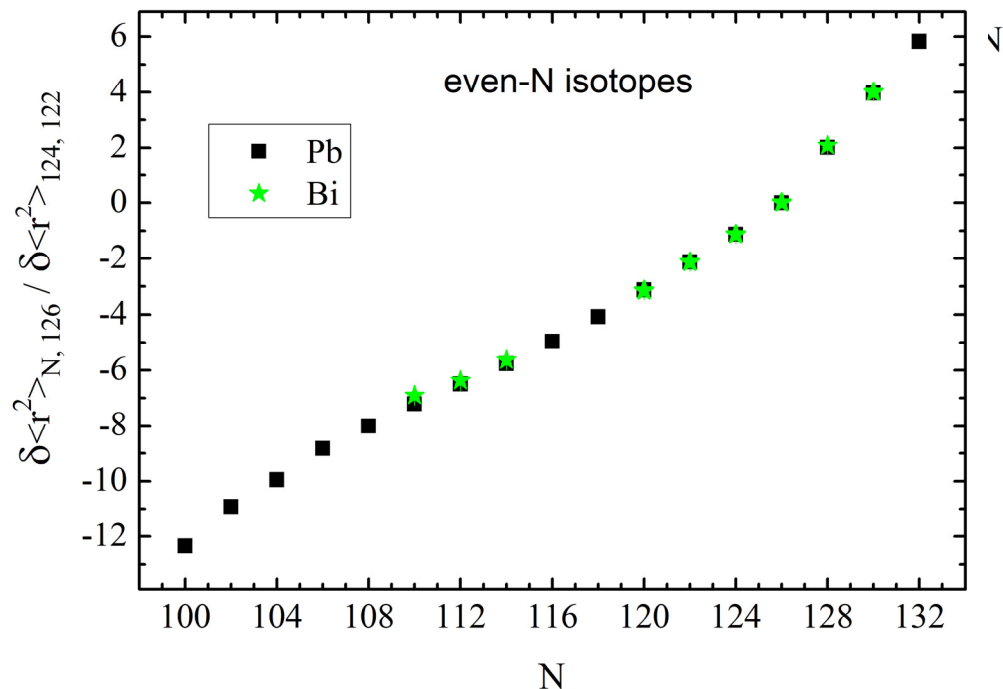
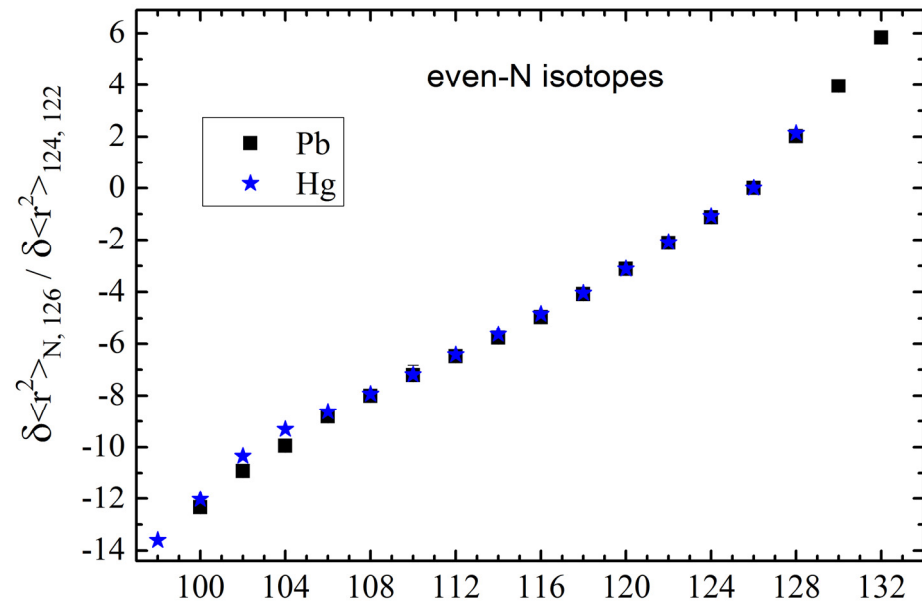
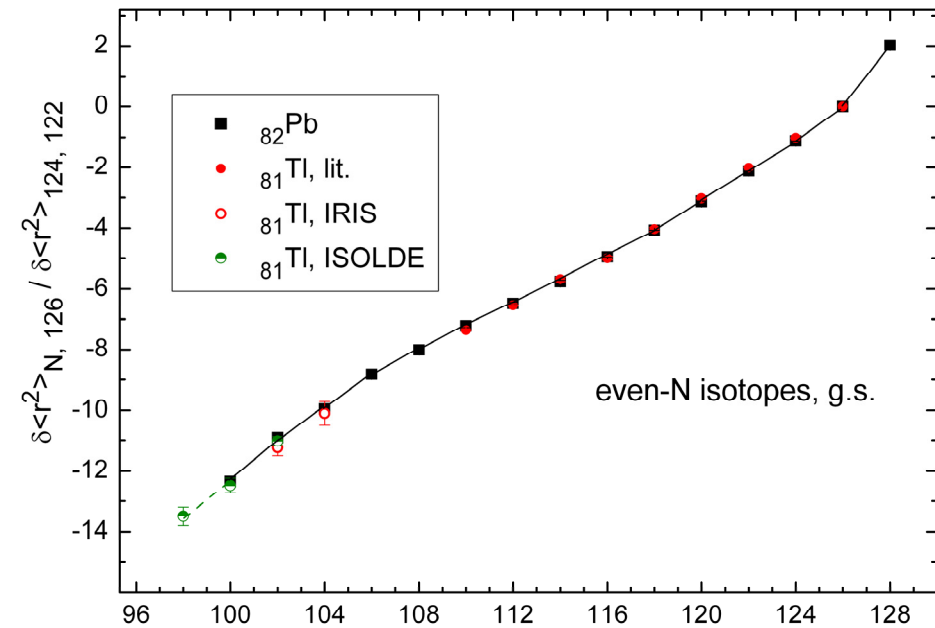
Due to limited resolution, for the spin determination the “integration method” was used.

The ratio of areas under each resolved peak:

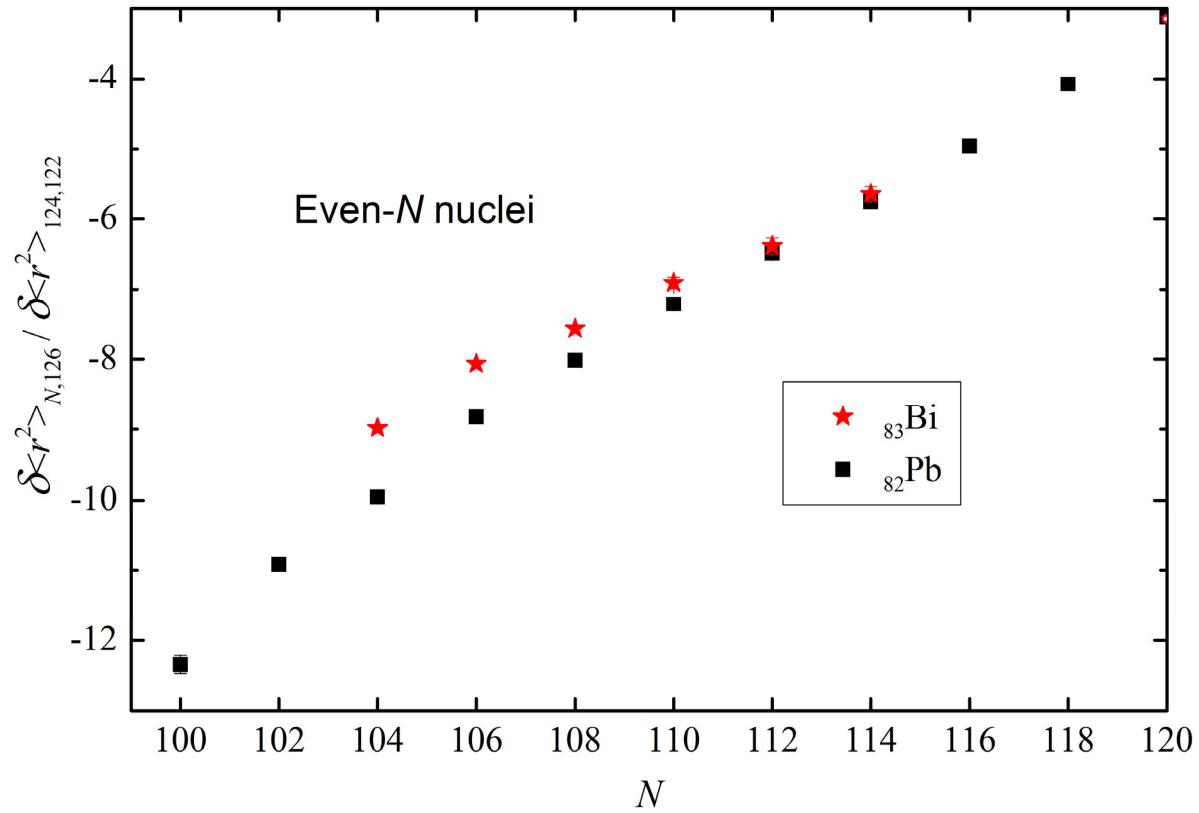
$$r_{\text{theor}} = (I+1)/I \cdot r_{\text{theor}} = 2 \text{ at } I = 1$$
$$r_{\text{theor}} = 1.5 \text{ at } I = 2$$
$$r_{\text{theor}} = 1.67 \text{ at } I = 3 \text{ etc.}$$

The weighted mean value, $r_{\text{expt}} = 2.00(12)$, for the six hfs spectra available for $^{188}\text{Bi}^g$ indicates a strong preference for an $I = 1$ assignment.

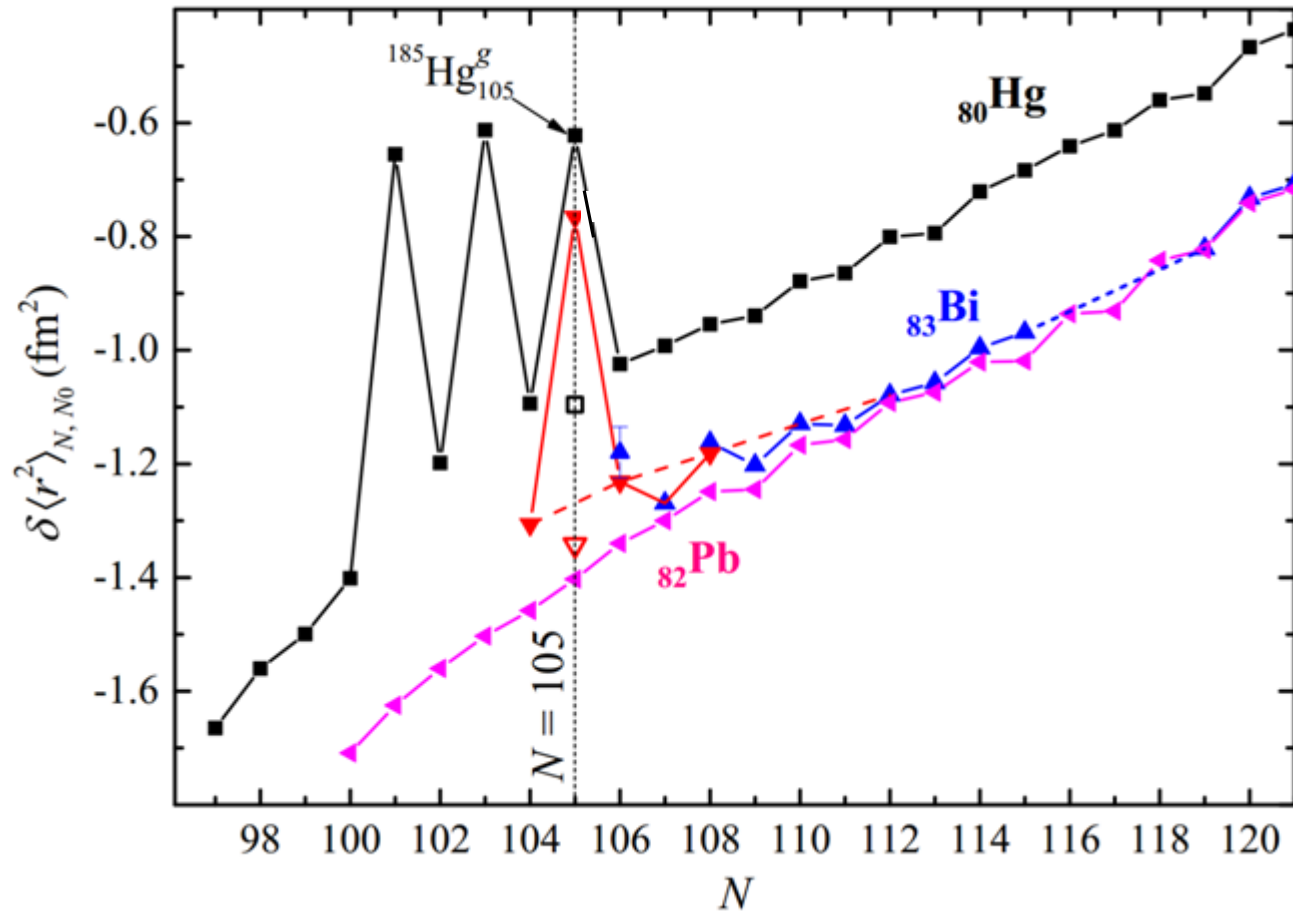
Relative radii: universal Pb trend



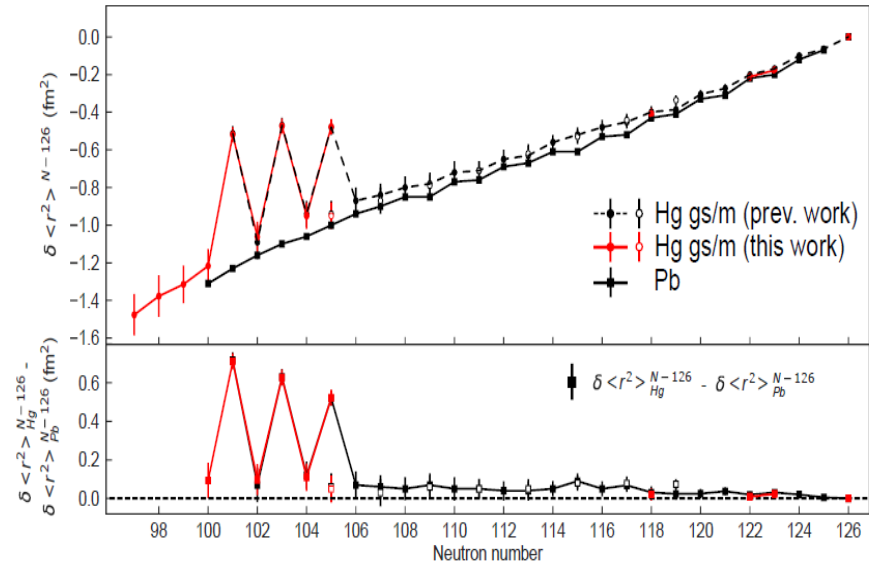
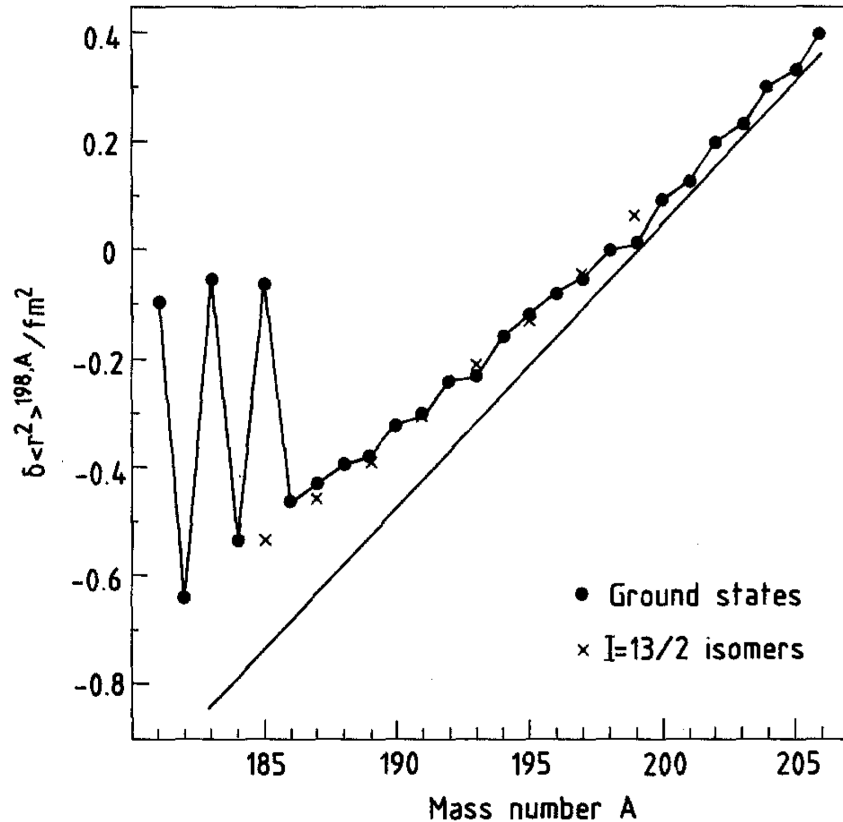
Relative radii: deviation from universal Pb trend



Shape staggering: Bi



Shape staggering: Hg



“One of the most remarkable discoveries in nuclear structure physics in the last 50 years”.
 K. Heyde and J. L. Wood, Phys. Scripta 91, 083008 (2016)

Shape staggering: $^{188}\text{Bi}g$, magnetic and quadrupole moments

$$\mu_{\text{exp}}(^{188}\text{Bi}; I = 1) = 0.994(21) \mu_N$$

Configuration with proper magnetic moment:

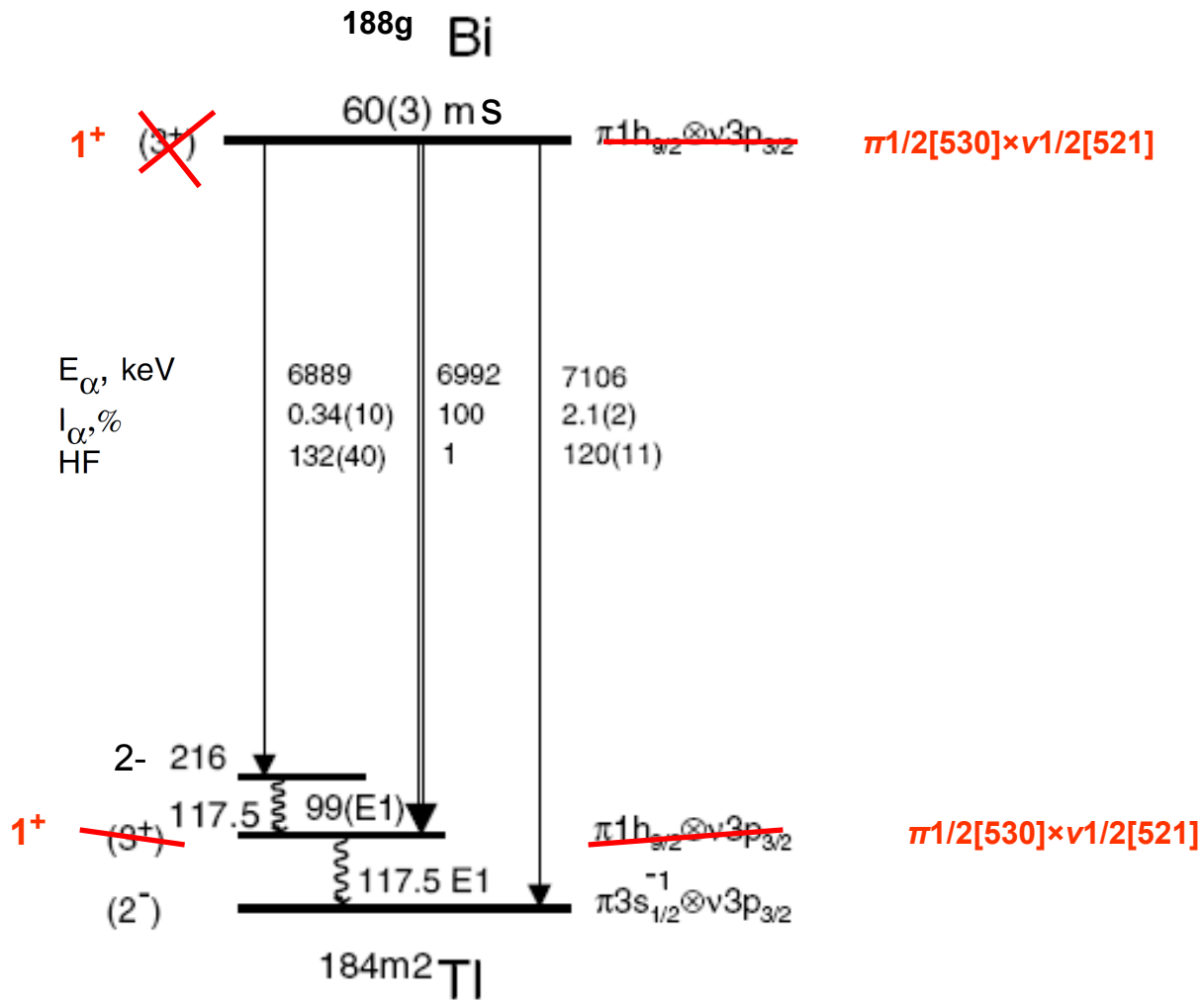
$$\mu_{\text{theor}}(\{\pi 1/2[530] \times \nu 1/2[521]\}_1) = 1.0(2) \mu_N$$

neutron orbital $1/2[521]$ is exactly the same that determines the strong prolate deformation in $^{181}, ^{183}, ^{185}\text{Hg}$; ^{185}Hg has the same neutron number ($N = 105$) as ^{188}Bi .

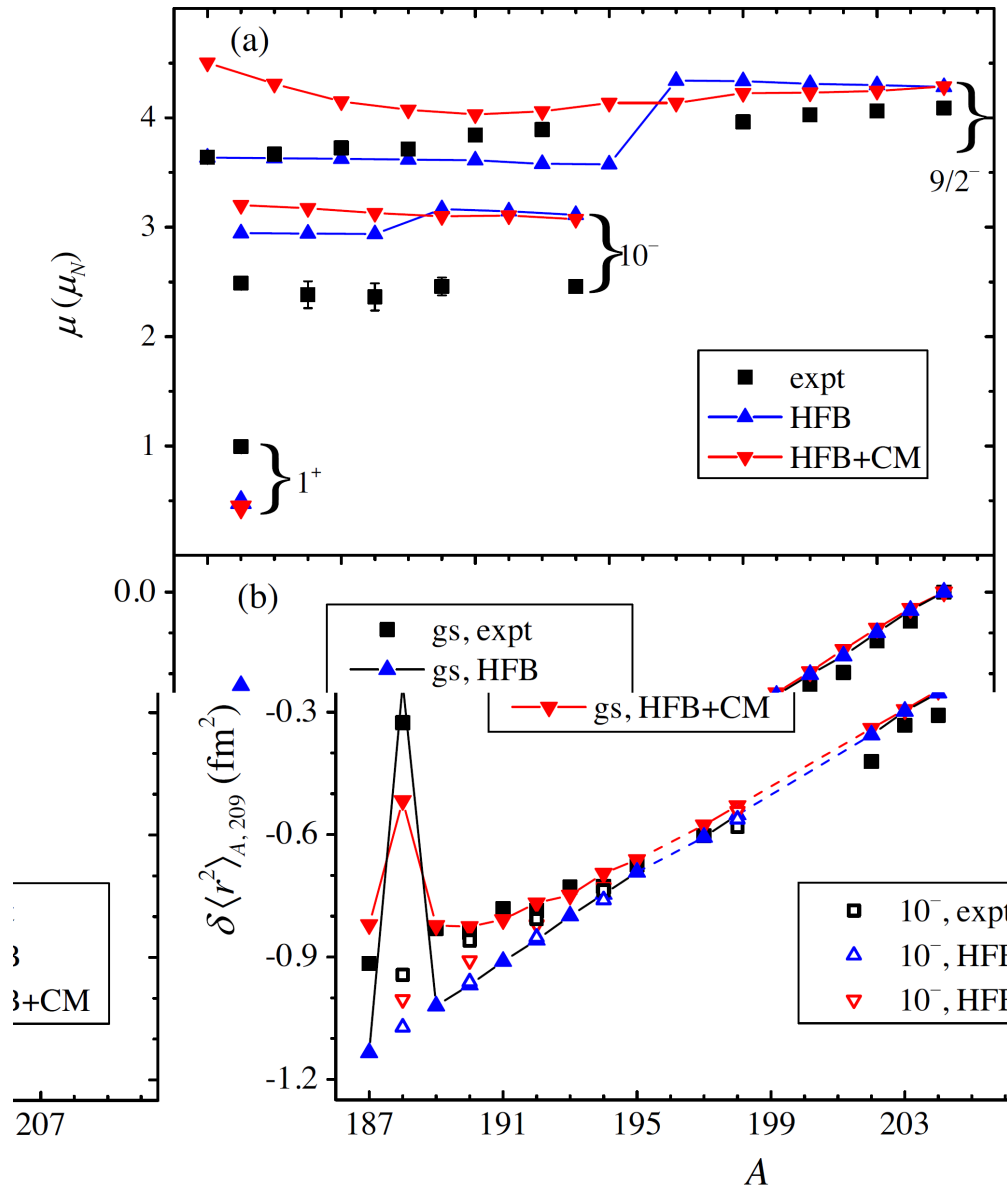
$$Q_{\text{exp}}(^{188}\text{Bi}g) = 0.85(37) \text{ b} \quad \Longrightarrow \quad \beta = +0.25(7)$$

$$Q_s = \frac{I \cdot (2I - 1)}{(I + 1) \cdot (2I + 3)} \cdot \frac{3}{\sqrt{5\pi}} \cdot Z \cdot R_0^2 \cdot \beta_Q \cdot \left(1 + \frac{1}{7} \cdot \sqrt{\frac{20}{\pi}} \cdot \beta_Q + \dots\right)$$

^{188}Bi : α decay



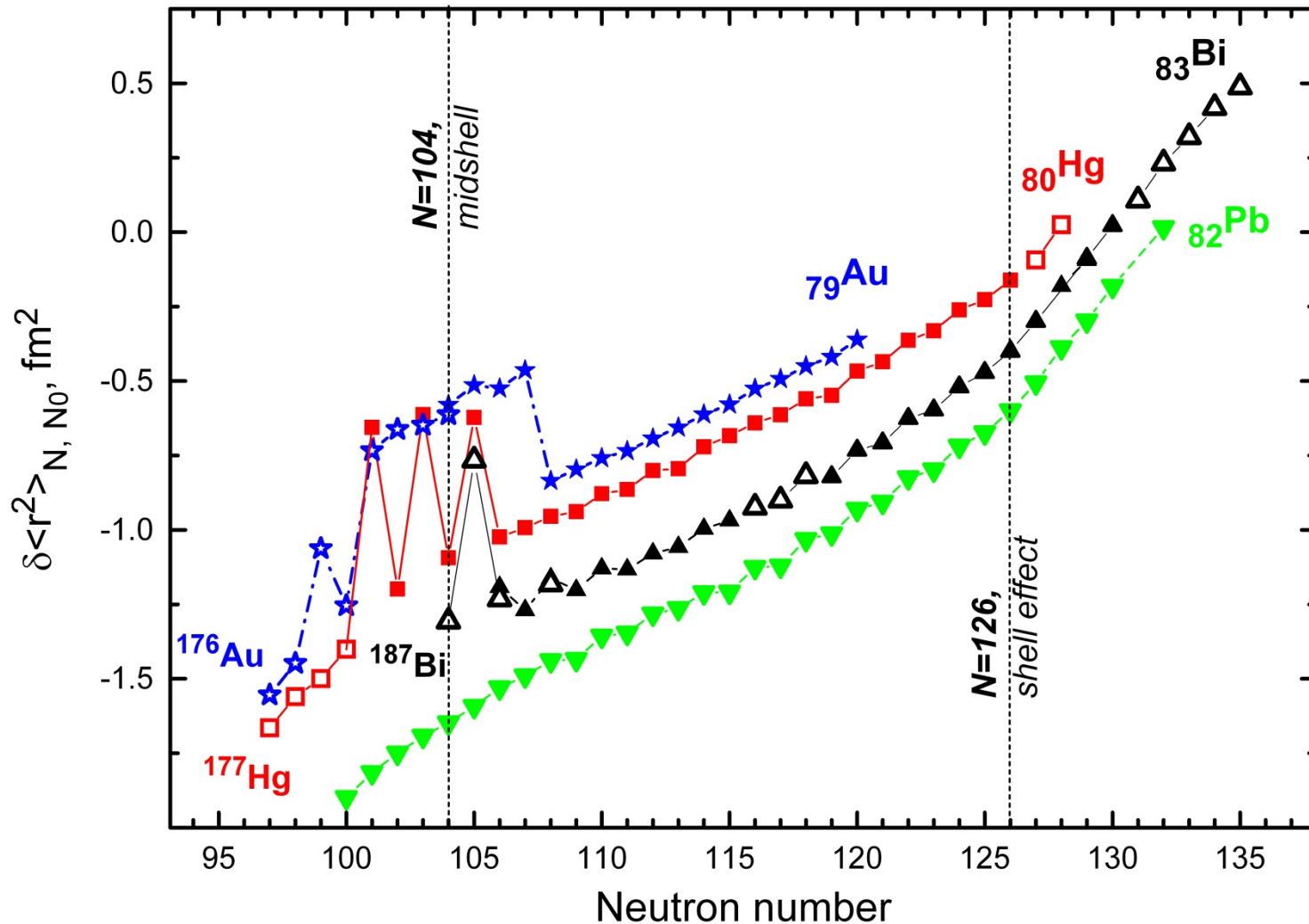
Shape staggering and OES in Bi: theory



Phenomenological model of configuration mixing:

$$\langle \mathcal{O} \rangle = \frac{\int_q \mathcal{O} \exp(-E/T) dq}{\int_q \exp(-E/T) dq}$$

Different patterns of shape evolution near midshell



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Bismuth isotopes also alternate from spheres to rugby balls

The unusual nuclear physics phenomenon, first discovered at CERN's ISOLDE facility 50 years ago, had until now been seen only in mercury isotopes

10 NOVEMBER, 2021 | By Ana Lopes



The ultra-rapid laser set up used by the ISOLDE team to study bismuth isotopes. Image: CERN.

Alternating from spheres to rugby balls is no longer the sole preserve of mercury isotopes, an international team at CERN's [ISOLDE](#) facility reports in a [paper](#) published in *Physical Review Letters*.

Isotopes are forms of a chemical element that have the same number of protons in their atomic nuclei but a different number of neutrons.

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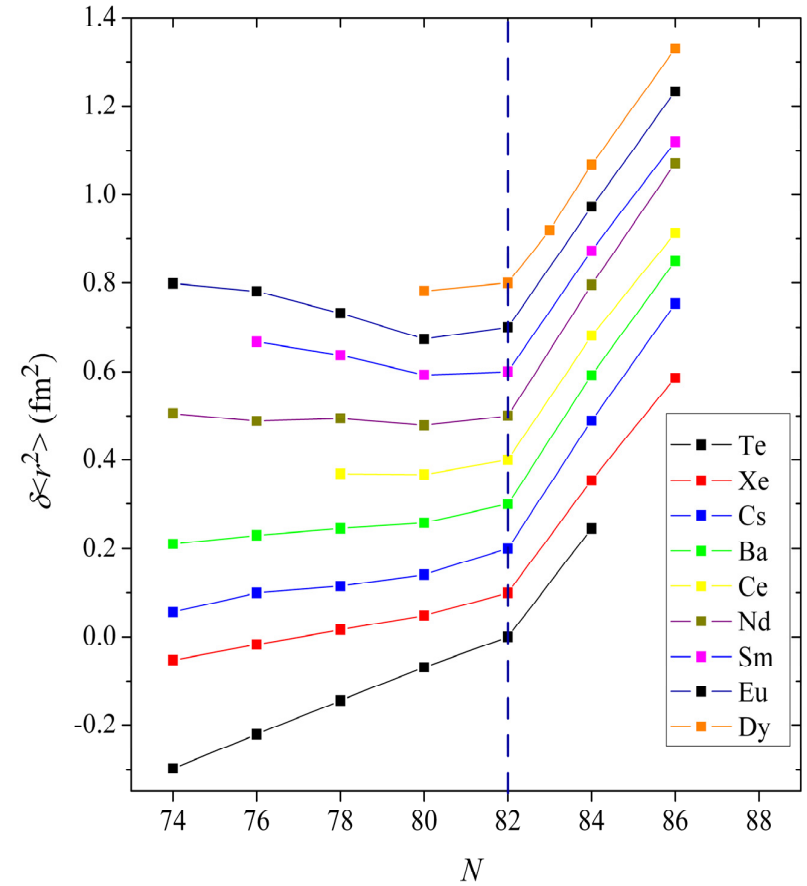
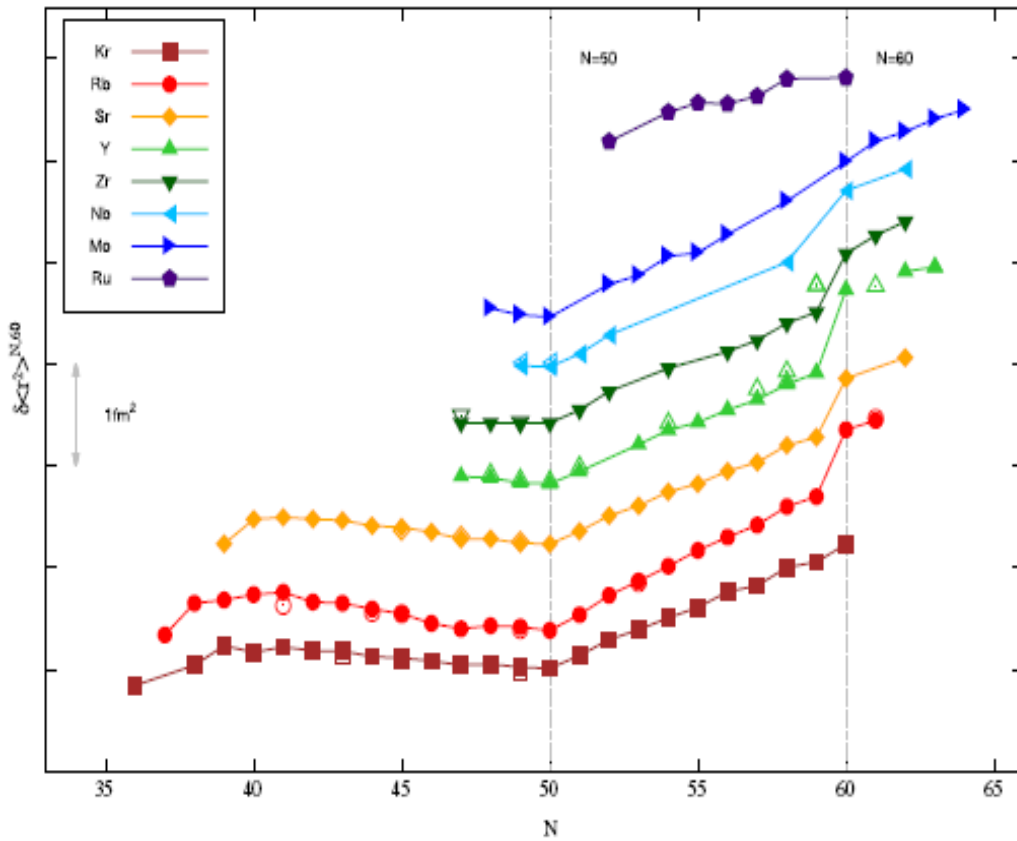
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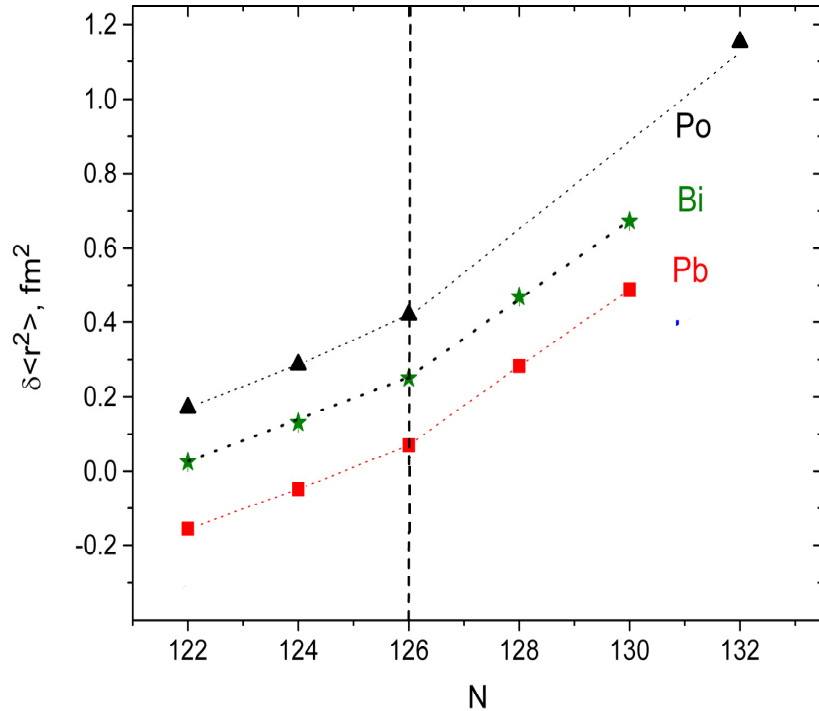
Barzakh A., ... Fedorov D. V., ... Molkanov P., ... Pantelev V., ... Skripnikov L. V., ... Zaitsevskii A. V.,
Large Shape Staggering in Neutron-Deficient Bi Isotopes, Phys. Rev. Lett. **127**, 192501 (2021).

Shell effect in radii at $N = 50, 82$

P. Campbell et al., Progress in Particle and Nuclear Physics 86 (2016) 127–180

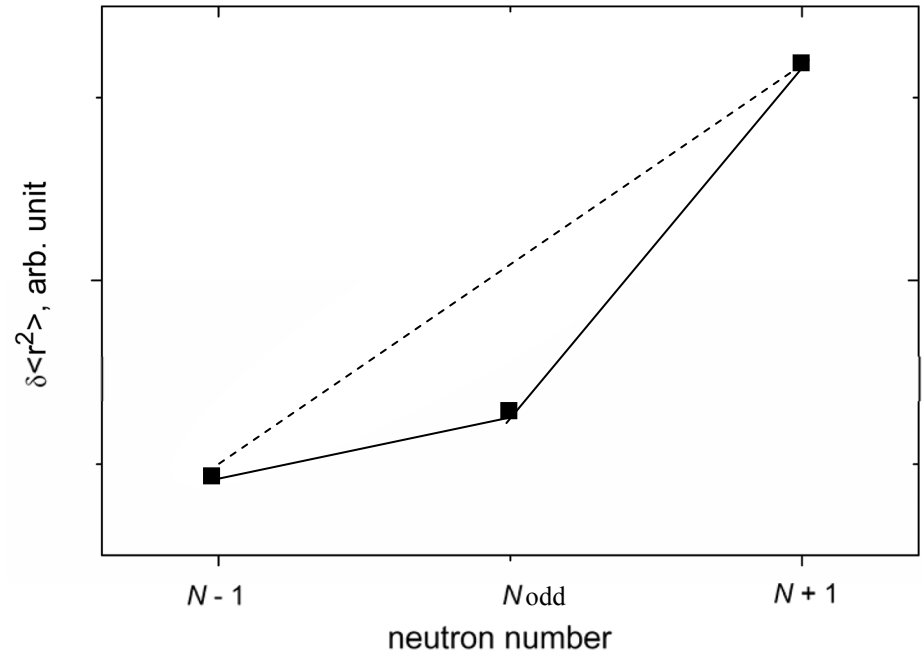


Shell effect in radii at $N = 126$ and odd-even staggering

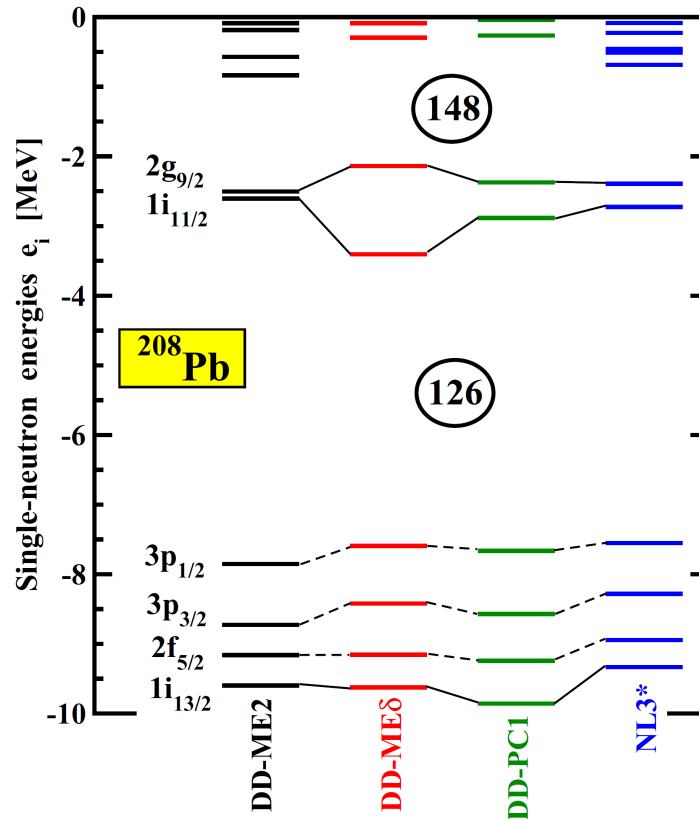
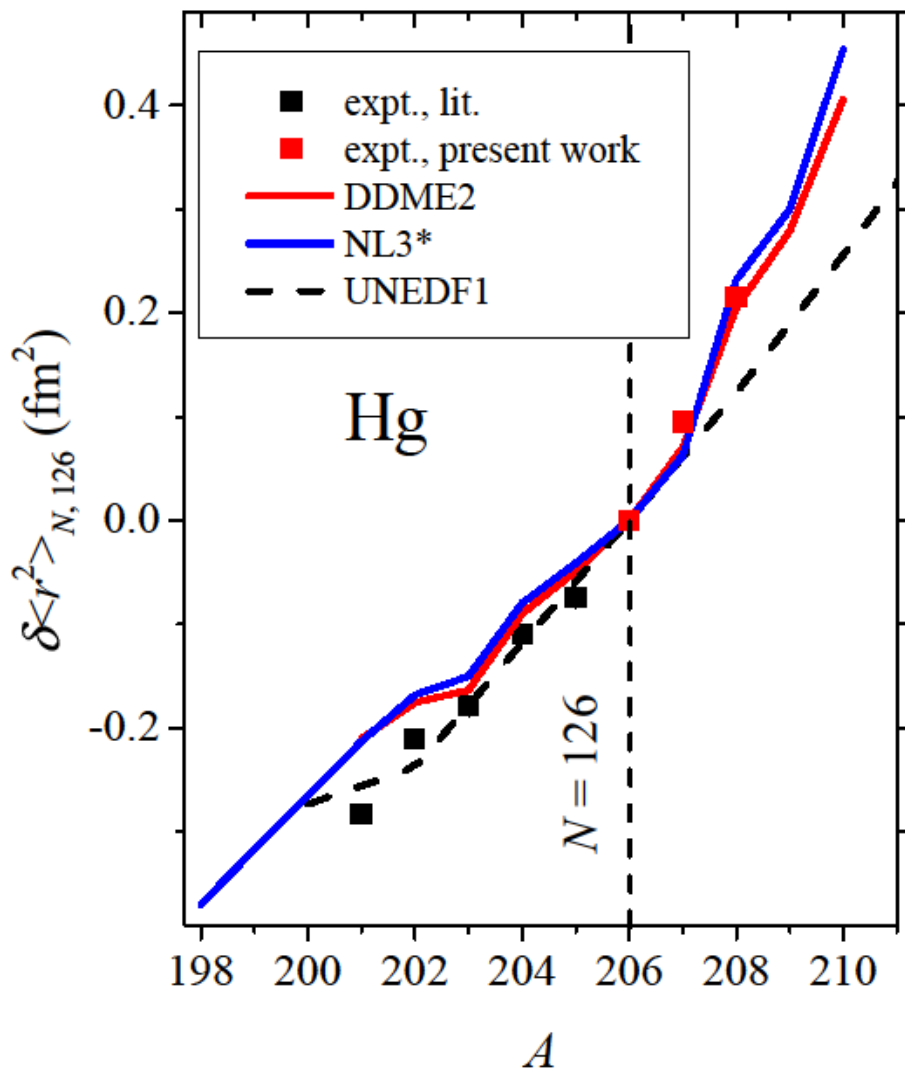


Odd-even staggering (OES) —
another universal feature

The shell effect in the changes of the nuclear mean-square charge radius $\delta\langle r^2 \rangle$ — the kink in its isotopic trend at the magic neutron numbers — was found to be an universal feature of the $\delta\langle r^2 \rangle$ behavior



Shell effect in radii: comparison with RMF



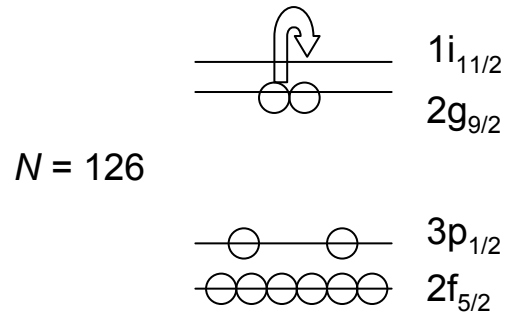
$$I(^{207}\text{Hg}_{127}) = 9/2$$

$$\mu(^{207}\text{Hg}_{127}) = -1.37(2) \mu_N$$

$$\mu_{\text{sp}}(\text{vg}_{9/2}) = -1.3 \mu_N$$

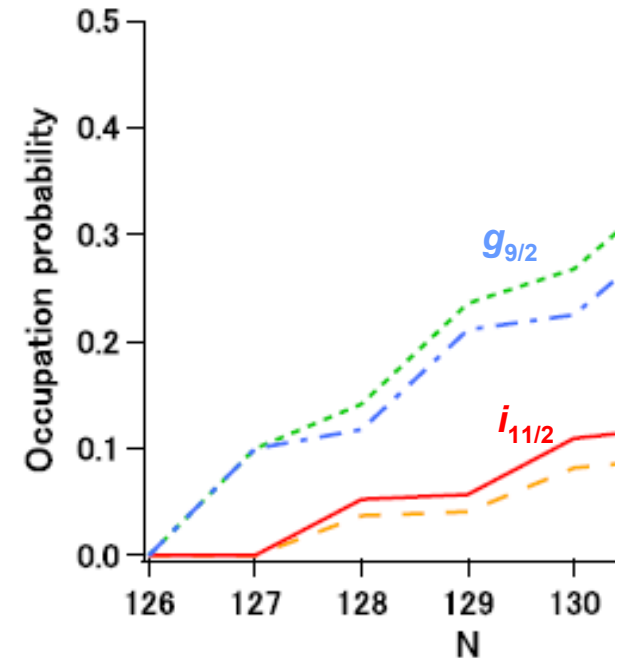
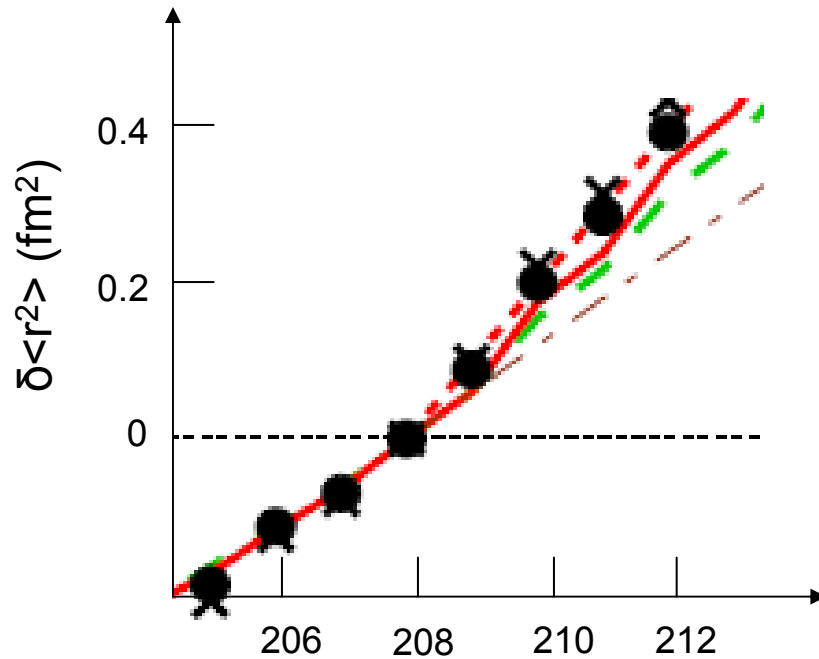
Shell effect in radii: theory

P. M. Goddard, P. D. Stevenson, and A. Rios, Phys. Rev. Lett. **110**, 032503 (2013):
The reproduction of the isotope shift in lead is determined by
the occupation of the $1i_{1/2}$ neutron orbital.



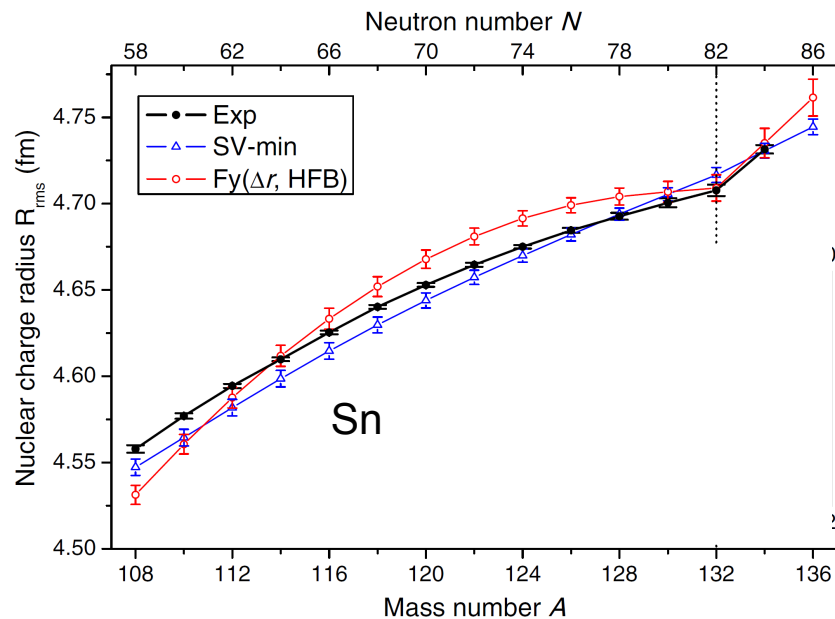
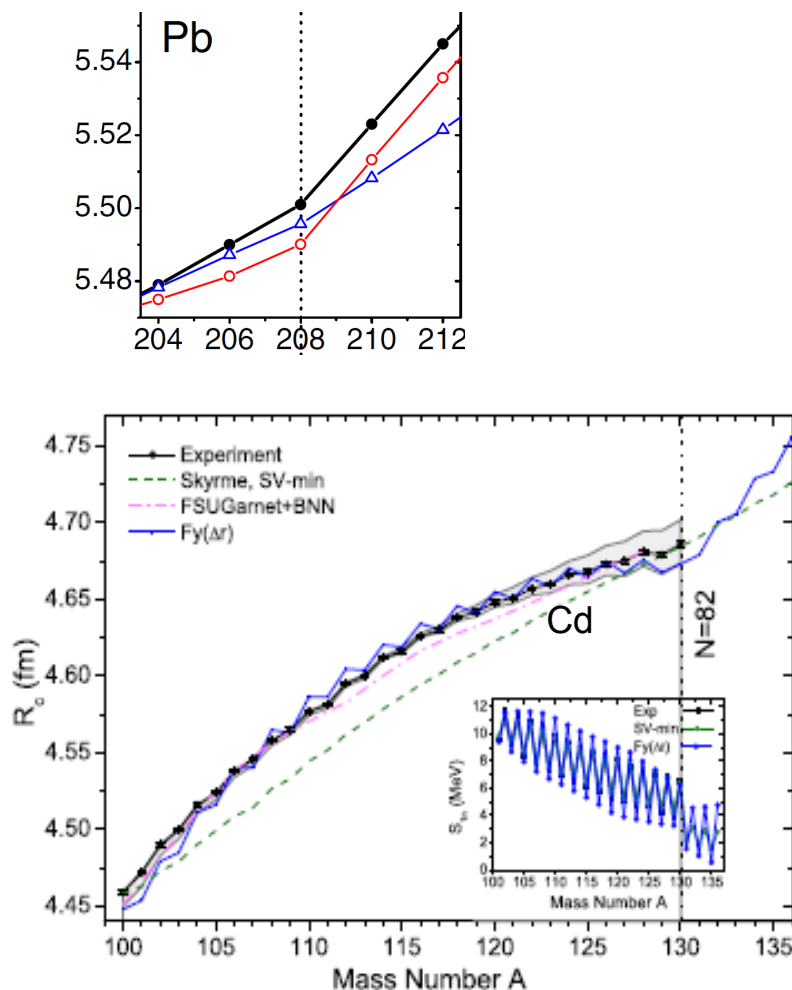
The increase of the $\nu i_{1/2}$ -orbital occupancy in the RMF and modified HF approaches was shown to be connected with the decrease of the energy splitting between the $\nu g_{9/2}$ and $\nu i_{1/2}$ levels in contradiction with the experimental evidences

Shell effect in radii: comparison with NRMF



H. Nakada succeeded in reproducing the $\delta\langle r^2 \rangle$ behavior for the Hg nuclei in the HFB calculations with density-dependent LS interaction based on 3N forces from effective field theory. The isotope shifts were described without the $ng_{9/2}$ - $ni_{11/2}$ degeneracy.

Shell effect in radii: Fayans functional



M. Hammen et al., Phys. Rev. Lett. 121, 102501 (2018)
 C. Gorges et al., Phys. Rev. Lett. 122, 192502 (2019)

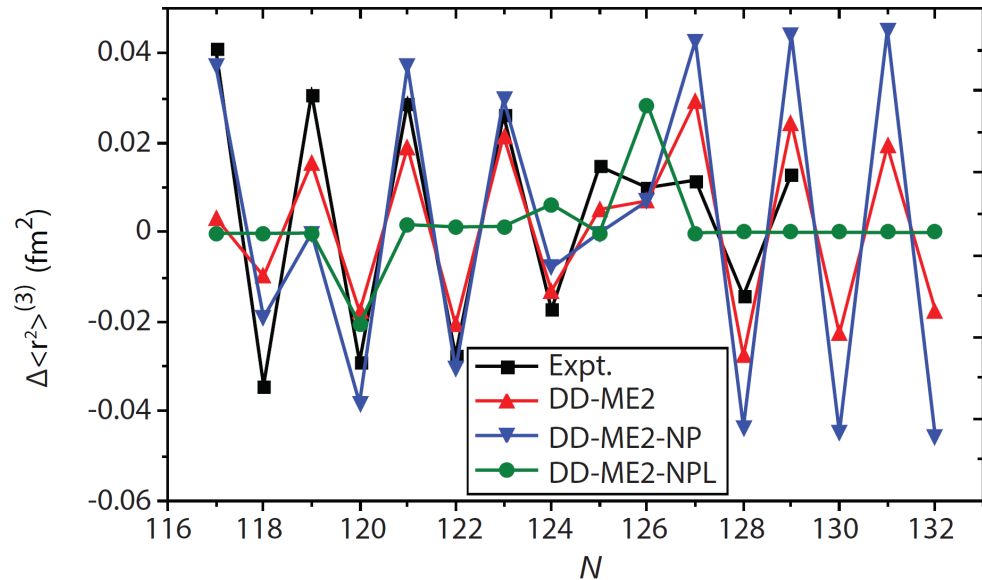
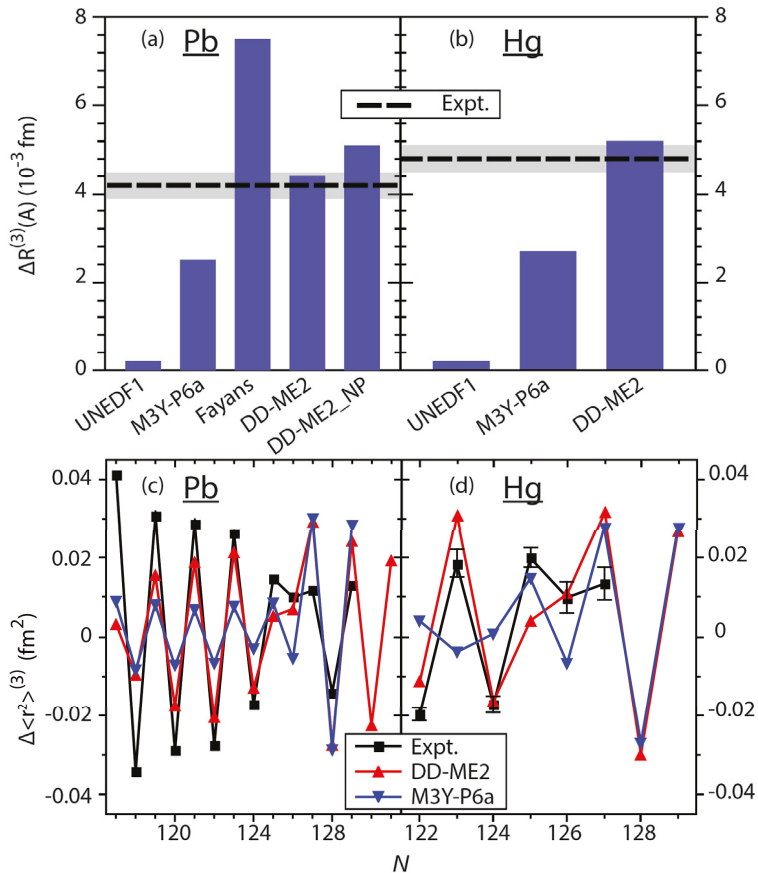
Energy density functional method based on the theory of finite Fermi systems. Fayans functional involves gradient terms in surface and pairing energies. They are responsible for the kink description.
 S. A. Fayans, S. V. Tolokonnikov, E. L. Trykov, D. Zawischa, Nucl. Phys. A **676**, 49 (2000)
 P.-G. Reinhard and W. Nazarewicz, Phys. Rev. C **95**, 064328 (2017)

Shell effect in radii: comparison of theoretical approaches

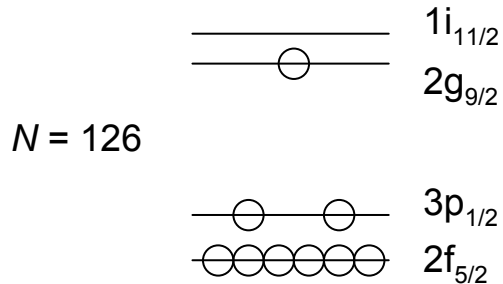
OES indicator: $\Delta \langle r^2 \rangle^{(3)}(N) = \frac{1}{2} [\langle r^2 \rangle(N-1) + \langle r^2 \rangle(N+1) - 2\langle r^2 \rangle(N)]$

kink indicator: $\Delta R^{(3)}(A) = \frac{1}{2} [R(A-2) + R(A+2) - 2R(A)]$,

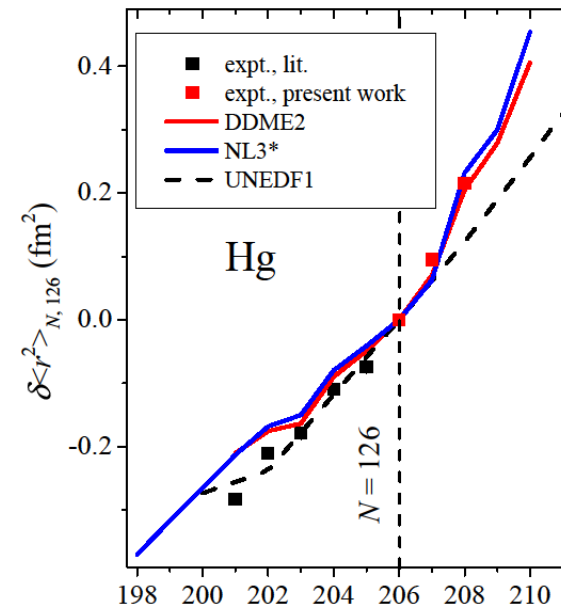
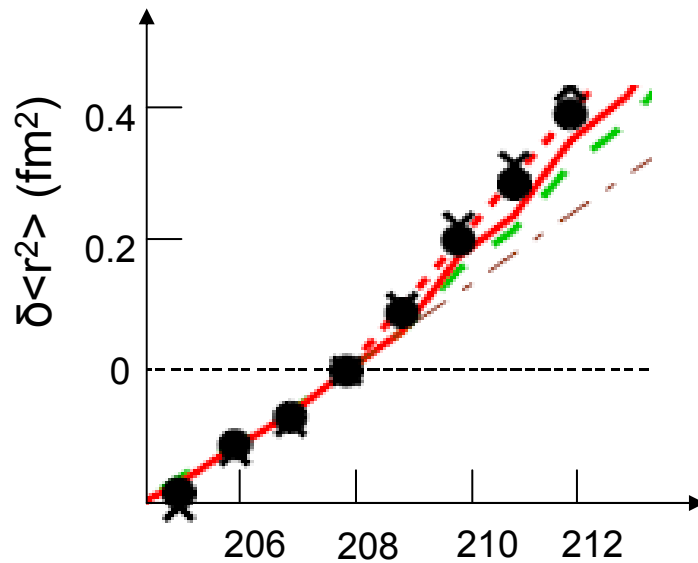
where $R(A) = \langle r^2 \rangle^{1/2}(A)$



Shell effect in radii: odd- N nuclei



For $N = 127$ the occupancy of the $n1i_{11/2}$ state is equal to zero (pairing is absent). Correspondingly, all models with the kink explanation by the increase of this occupancy predict $\xi_{\text{odd}} \sim 1$ in contradiction with experiment



1. Ядерная спектроскопия $^{176,177,179}\text{Au}$: схемы распада, времена жизни, факторы задержки альфа распада, схемы уровней дочерних ядер и т. д.
2. α - и β -распад $^{183}\text{Tl}^m$: деформированные возбужденные состояния в сферическом ^{179}Au .
3. Ядерная спектроскопия ^{214}Bi : обнаружен новый долгоживущий изомер, существенно расширена схема уровней дочернего ^{214}Po , проверка продвинутых shell-model расчетов.
4. Времена жизни возбужденных уровней ^{214}Po
5. Анализ выходов изотопов золота из урановой мишени на установке ISOLDE.
6. Атомные расчеты аномалии стс в золоте и градиента электрического поля электронов на ядре висмута (для квадрупольного момента ядра)
7. Изомерно-селективная ядерная спектроскопия вблизи дважды магического ^{132}Sn
8. Запаздывающее деление ^{178}Au

1. Barzakh A., ... Fedorov D.V., ... Molkanov P., ... Panteleev V., ... Skripnikov L.V., ... Zaitsevskii A.V., *Large Shape Staggering in Neutron-Deficient Bi Isotopes*, Phys. Rev. Lett. **127**, 192501 (2021).
2. T. Day Goodacre, A. V. Afanasjev, A. E. Barzakh, ... D. V. Fedorov, ... P. L. Molkanov, ... M. D. Seliverstov, et al., *Laser Spectroscopy of Neutron-Rich $^{207,208}\text{Hg}$ Isotopes: Illuminating the Kink and Odd-Even Staggering in Charge Radii across the $N = 126$ Shell Closure*, Phys. Rev. Lett. **126**, 032502 (2021).
3. T. Day Goodacre, A. V. Afanasjev, A. E. Barzakh, ... D. V. Fedorov, ... P. L. Molkanov, ... M. D. Seliverstov, et al. *Charge radii, moments, and masses of mercury isotopes across the $N = 126$ shell closure*, Phys. Rev. C **104**, 054322 (2021).
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6. R. D. Harding, A. N. Andreyev, A. E. Barzakh, ... D. V. Fedorov, ... P. L. Molkanov, ... M. D. Seliverstov et al., *Laser-assisted nuclear decay spectroscopy of $^{176,177,179}\text{Au}$* , Phys. Rev. C **104**, 024326 (2021).
7. B. Andel, ... A. Barzakh, ... D. V. Fedorov, ... P. Molkanov, ... M. D. Seliverstov, ... (IDS Collaboration), *New β -decaying state in ^{214}Bi* , Phys. Rev. C **104**, 054301 (2021).