

Is the $\mathcal{LHCb} \mathcal{P}_c(4312)^+$ Plausible in the GlueX $\gamma p \rightarrow J/\psi p$ Total Cross Sections?



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Joint work with Bill Briscoe
Eugene Chudakov
Ilya Larin
Lubomir Pentchev
Axel Schmidt
Ron Workman

- J/ψ photoproduction is sample of hard process corresponding to relatively large scale.
- Quality of novel GlueX J/ψ photoproduction data & proximity of data to energy threshold, gives access to variety of interesting physics aspects.
- Examples are *trace anomaly*, estimation of J/ψ - N scattering length, $5q$ exotics, & so on.

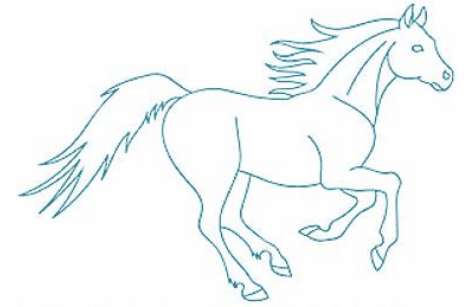


IIS, W.J. Briscoe, E. Chudakov, I. Larin, L. Pentchev, A. Schmidt, R.L. Workman,
Phys Rev C **108**, 015202 (2023)

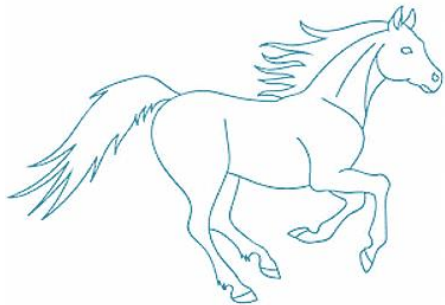
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Outline

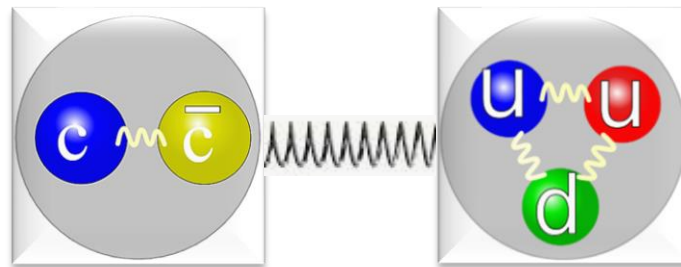


- *Narrow pentaquarks from $LHCb$ - Big question.*
- *$GlueX$ J/ψ photoproduction in 2019 & 2023.*
- *Intermission: Quantum interference of particles & resonances.*
- *Alternative solution for $GlueX$ J/ψ data.*
- *Cusp effect.*
- *Vector meson – nucleon scattering length.*
- *$J\Phi AC$ for J/ψ photoproduction.*
- *$SOLID$ for more J/ψ photoproduction.*
- *$JLab$ at High Energies.*
- *$J-PARC$ for J/ψ .*
- *Summary.*



Narrow Pentaquarks from LHCb

- Big Question



Narrow Pentaquarks from $\Lambda_b \rightarrow J/\psi p K^-$

- QCD gives rise to *hadron spectrum*.

Volume 8, number 3 PHYSICS LETTERS 1 February 1964

A SCHEMATIC MODEL OF BARYONS AND MESONS CI=4032

M. GELL-MANN

If we assume that the strong interactions of baryons and mesons are correctly described in terms of the broken "eightfold way" Baryons can now be constructed from quarks by using the combinations (qqq) , $(qqq\bar{q})$, etc., while mesons are made out of $(q\bar{q})$, $(qq\bar{q}\bar{q})$, etc.

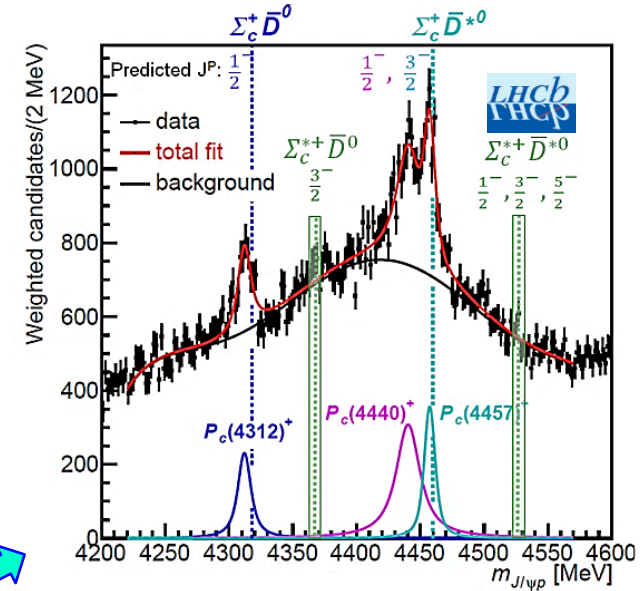
- Many $\bar{q}q$ & qqq states have been *observed*.

220 & 100.

- $\bar{q}q\bar{q}q$, $qqq\bar{q}q$, ... are *not forbidden* or *we do not know it yet*.

• claims *evidence* for *four hidden-charm $qqq\bar{q}q$ states* near *open-charm* decay thresholds for $\Sigma_c^+ \bar{D}^0$ & $\Sigma_c^+ \bar{D}^{*0}$ in $\Lambda_b \rightarrow J/\psi p K^-$ decays.

State	M (MeV)	$\Gamma [P_c \rightarrow J/\psi + p]$ (MeV)	Significance
$P_c(4312)^+$	$4311.9 \pm 0.7_{-0.6}^{+6.8}$	$9.8 \pm 2.7_{-4.5}^{+3.7}$	7.3σ
$P_c(4337)^+$	4337_{-4-2}^{+7+2}	29_{-12-14}^{+26+13}	$3.1 - 3.7 \sigma$
$P_c(4440)^+$	$4440.3 \pm 1.3_{-4.7}^{+4.1}$	$20.6 \pm 4.9_{-10.1}^{+8.7}$	5.4σ
$P_c(4457)^+$	$4457.3 \pm 0.6_{-1.7}^{+4.1}$	$6.4 \pm 2.0_{-1.9}^{+5.7}$	5.4σ



R. Aaij *et al*, Phys Rev Lett **122**, 222001 (2019)

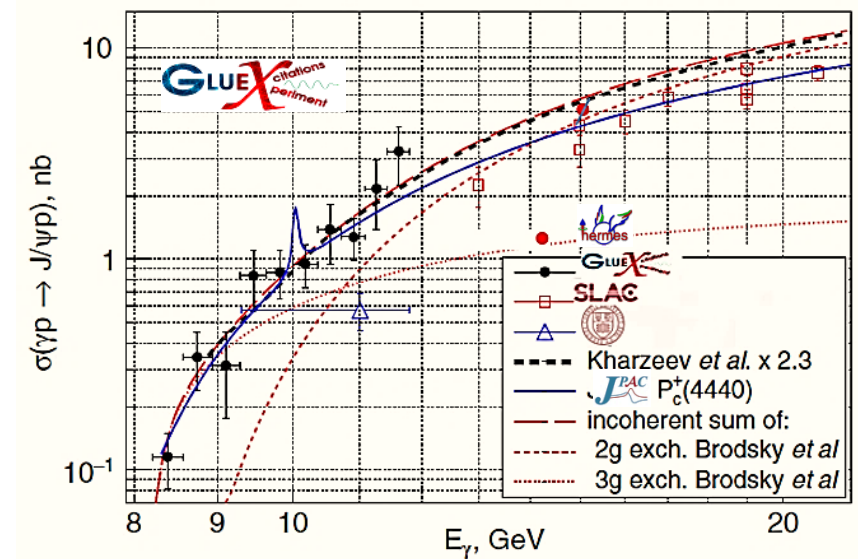
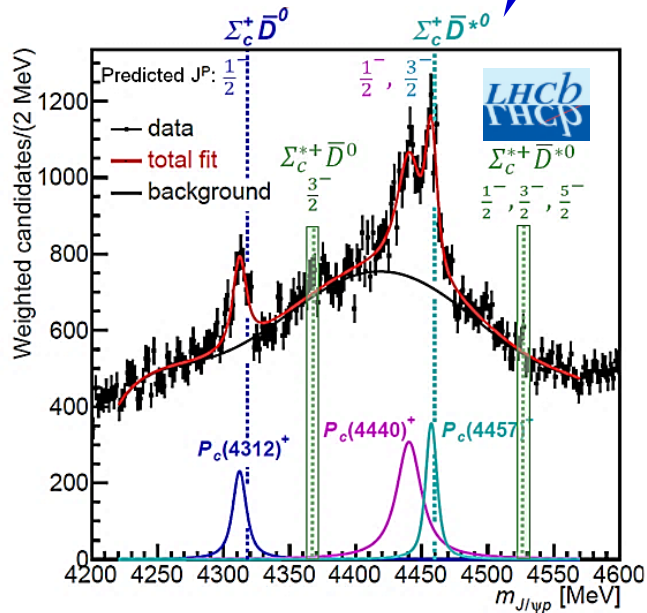
631

- **Bump hunting:**
 - no *quantum numbers*
 - no *pole positions*

Bump Hunting



- Why **LHCb** has chance *to see hidden-charm pentaquark states* while it *is elusive* in **GLUEX** ~~experiment~~?
- There is no confirmation of this **LHCb** observation **PDG**.

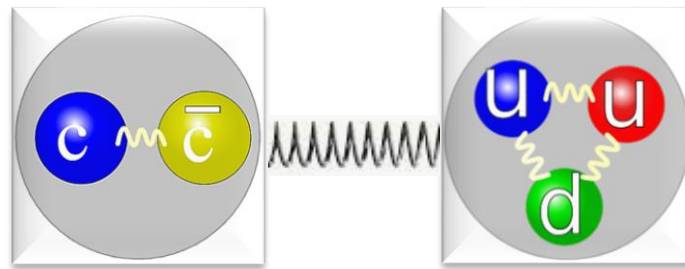


“Not every *bump* is a *resonance*, not every *resonance* is a *bump*”
 R. Gordon Moorhouse (1960s)



GlueX J/ψ Photoproduction

in 2019 & 2023



How Bump Hunting works in 2019 data?

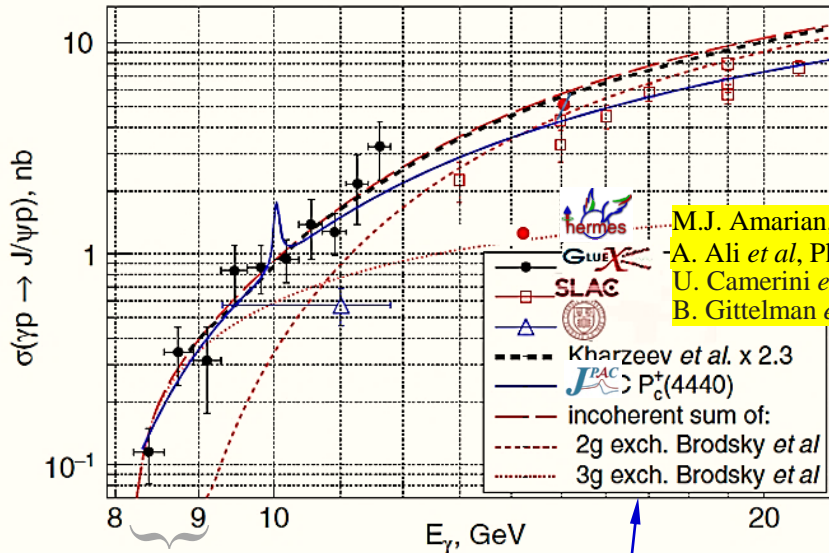
Editors' Suggestion

A. Ali *et al*, Phys Rev Lett **123**, 072001 (2019)



CI=170

2016–2017 data: $469 \pm 22 \text{ } \gamma p \rightarrow J/\psi p \rightarrow e^+ e^- p$ & 68 pb^{-1}





M.J. Amarian, Few-Body Syst Suppl. **11**, 359 (1999)
 A. Ali *et al*, Phys Rev Lett **123**, 072001 (2019)
 U. Camerini *et al*, Phys Rev Lett **35**, 483 (1975)
 B. Gittelman *et al*, Phys Rev Lett **35**, 1616 (1975)

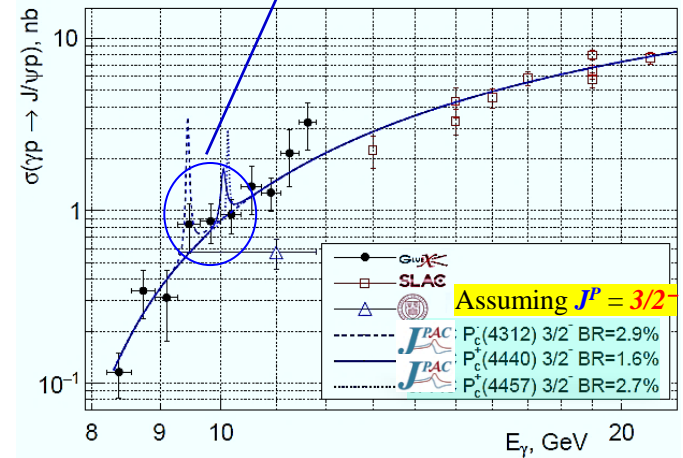
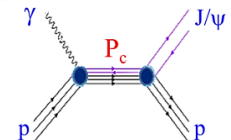
D. Kharzeev, H. Satz, A. Syamtomov, & G. Zinovjev, Nucl Phys A **661**, 568 (1999)
 J-PAC A.N. Hiller Blin *et al*, Phys Rev D **94**, 034002 (2016)
 S. Brodsky, E. Chudakov, P. Hoyer, & J.M. Laget, Phys Lett B **498**, 23 (2001)

• 3g works better than 2g @ threshold



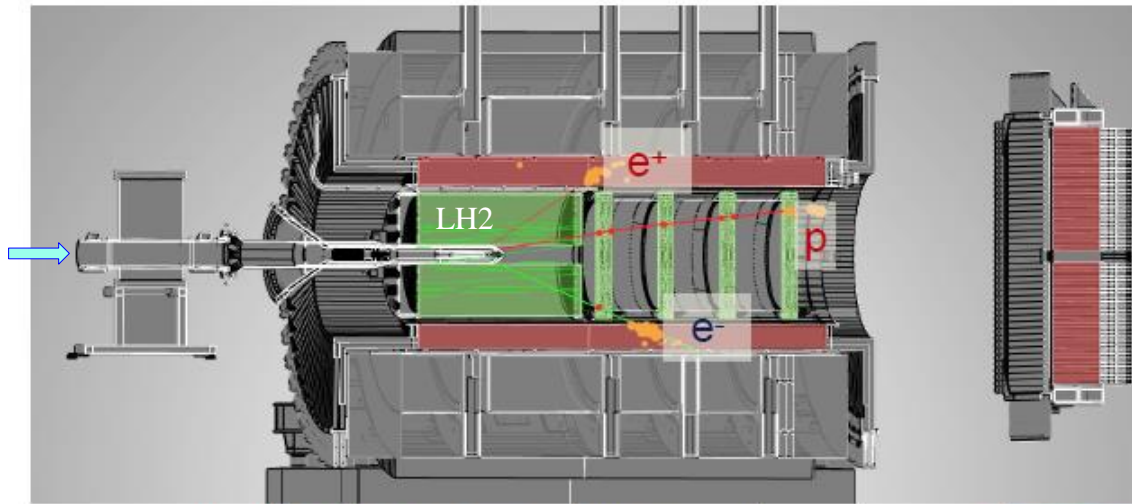
•  sees *no evidence* for  P_c s
 Upper limits @ 90% CL

State	Upper Limit
$P_c(4312)$	4.6 %
$P_c(4440)$	2.3 %
$P_c(4457)$	3.8 %



Recent (and future) Jefferson Lab results on threshold charmonium photoproduction

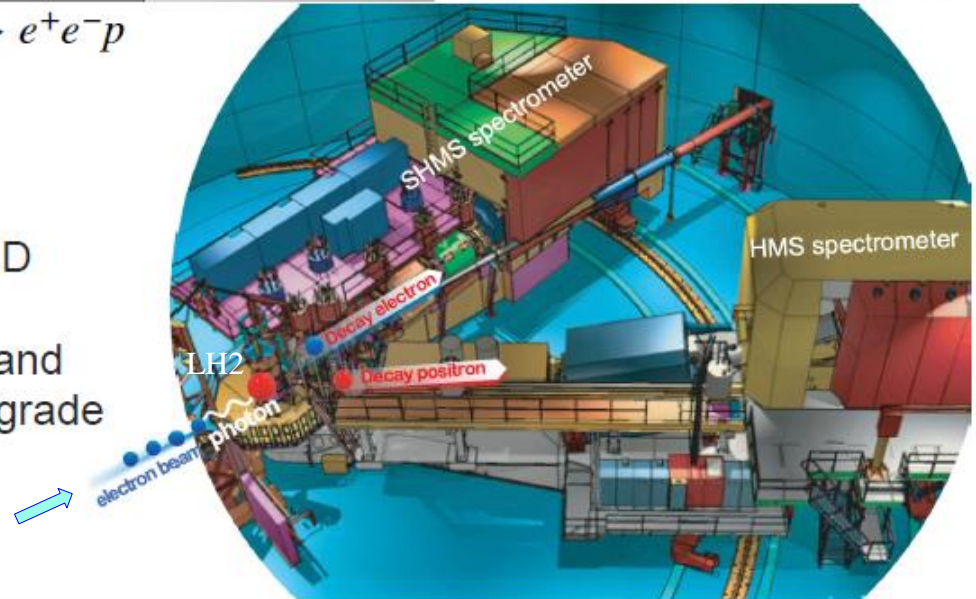
Lubomir Pentchev
GLUEX Collaboration



Hall D GLUEX $\gamma p \rightarrow J/\psi p \rightarrow e^+ e^- p$

Hall C 007 $(\gamma)p \rightarrow J/\psi(p) \rightarrow e^+ e^-(p)$

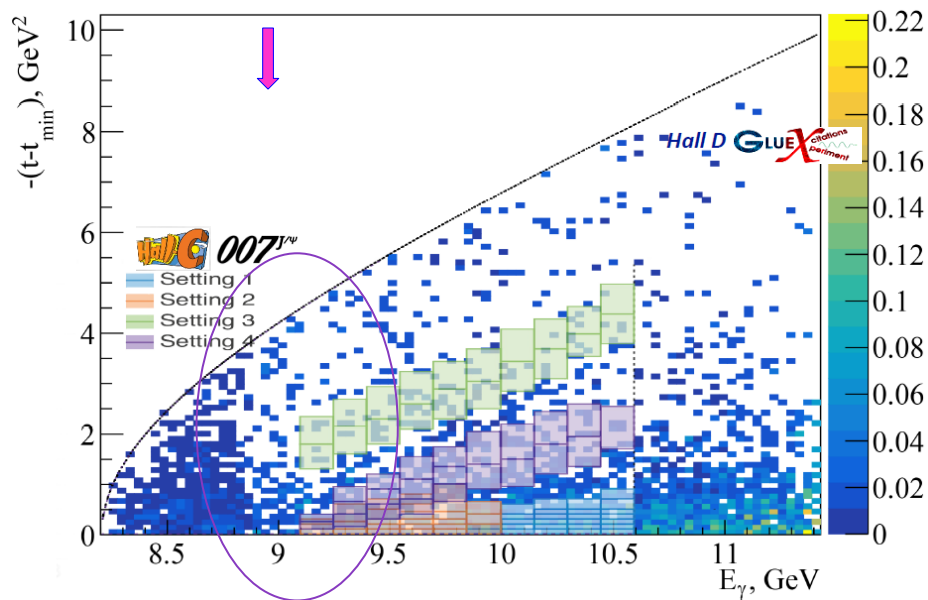
- Recent J/ψ results from Hall C & D
- ... and their interpretation
- Higher-mass charmonium states and prospects with CEBAF energy upgrade using GLUEX detector



Courtesy of Lubomir Pentchev, 2023

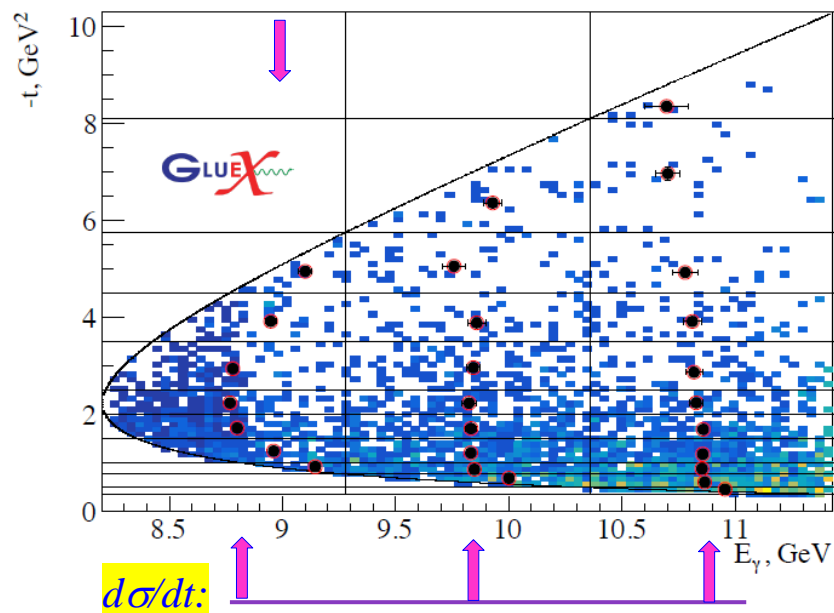


J/ψ Threshold Region Coverage



PDG $\text{BR}(J/\psi \rightarrow e^+e^-) = (5.971 \pm 0.032)\%$
 $\text{BR}(J/\psi \rightarrow \mu^+\mu^-) = (5.961 \pm 0.033)\%$

In progress



Courtesy of Lubomir Pentchev, 2023

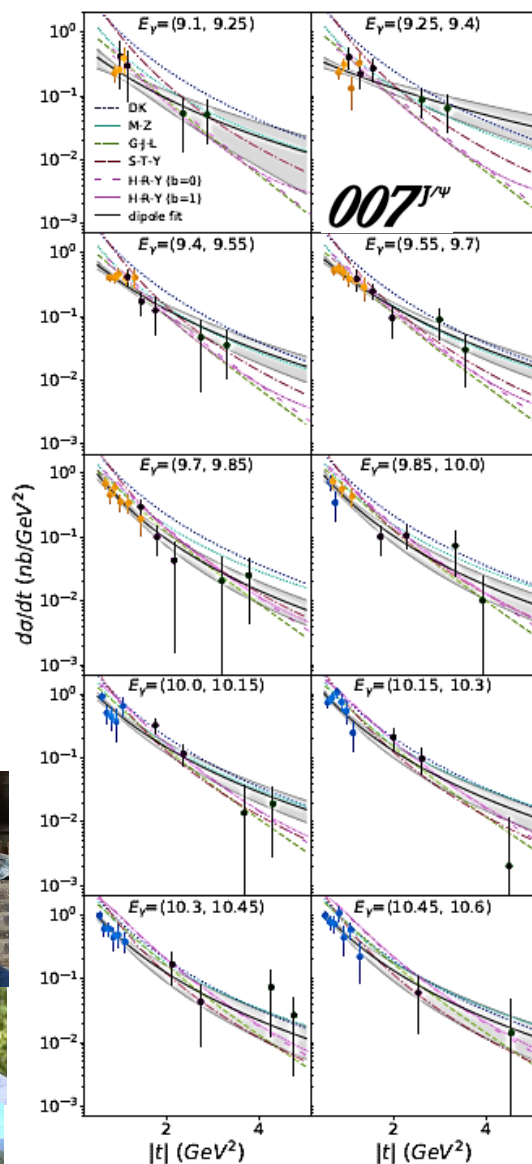
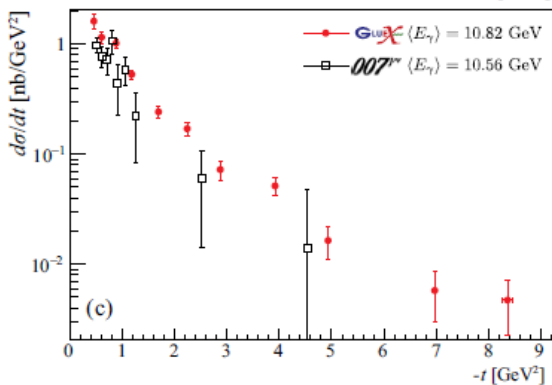
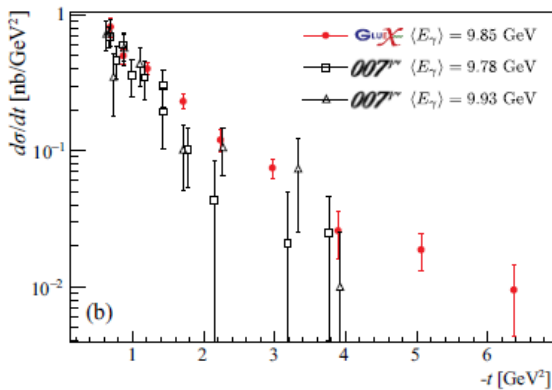
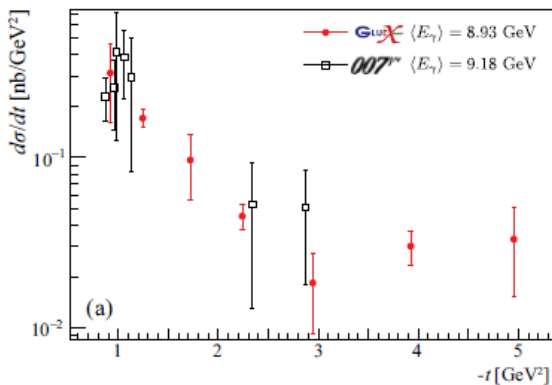
10/8/2023

Joint Seminar HEPD-THD, PNPI, Gatchina, Russia, October 2023

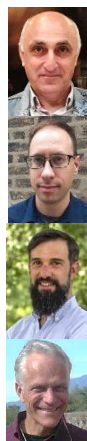
Igor Strakovsky 9



Differential Cross Sections from J/ψ by $007^{J/\psi}$ & GLUEX



- 10 energy bins in $007^{J/\psi}$.
- Results for 3 GLUEX energy bins compared to closes $007^{J/\psi}$ energies.
- Scale/Norm uncertainties: 20% in GLUEX & 4% in $007^{J/\psi}$ results.
- Good agreement within s.



S. Adhikari *et al*, Phys Rev C **108**, 025201 (2023)

CI=8



B. Duran *et al*, Nature **615**, 813 (2023)

CI=26



How Bump Hunting works in 2023 data?

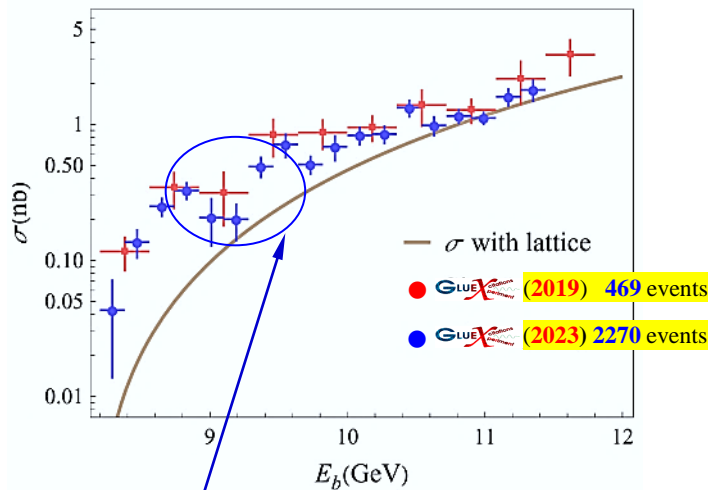
S. Adhikari *et al*, Phys Rev C **108**, 025201 (2023)




CI=10


Editors' Suggestion

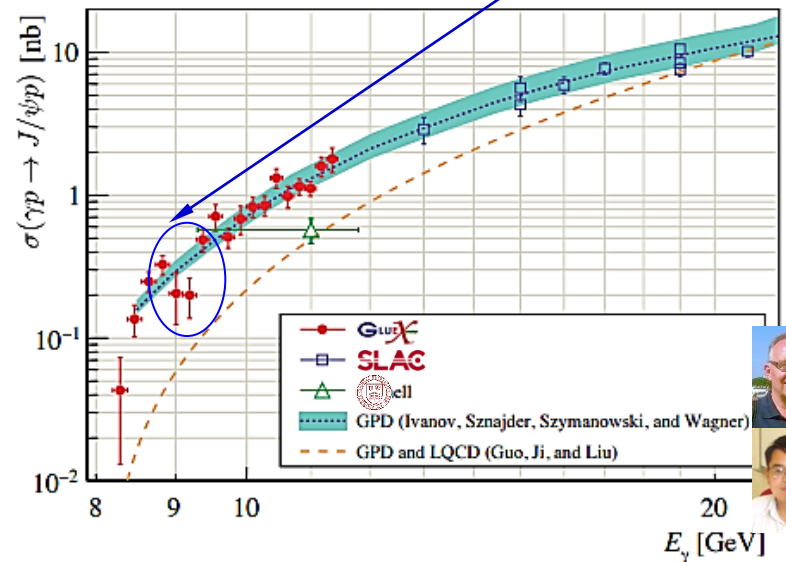
2016–2018 data: $2270 \pm 58 \gamma p \rightarrow J/\psi p \rightarrow e^+e^-p$ & 320 pb^{-1}



Y. Guo *et al*, Phys Rev D **108**, 034003 (2023)

- *Dip*, close to $E_\gamma = 9.2\text{-}9.3$ GeV, deepened with growing  statistics (by factor of 5).
- Data taken in 2020 will additionally increase statistics by factor of 2.

- If  treats two total Xsec points in vicinity of $W = 4.2\text{-}4.3$ GeV as potential *dip* (considering *stat* & *point-to-point syst* uncertainties), probability that they are not *stat fluctuation* from underlying smooth fit to observed Xsec corresponds to significance of **2.6 σ** .



Intermission

Quantum Interference of Particles & Resonances

IOP PUBLISHING

JOURNAL OF PHYSICS G: NUCLEAR AND PARTICLE PHYSICS

J. Phys. G: Nucl. Part. Phys. 37 (2010) 023001 (22pp)

doi:10.1088/0954-3899/37/2/023001

TOPICAL REVIEW

Quantum interference of particles and resonances

Ya Azimov

Petersburg Nuclear Physics Institute, Gatchina 188300, Russia



CI=18

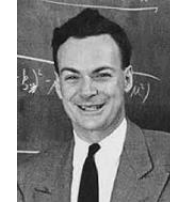


Interference - I



*When looking at **Maxwell** equations,
it is hard to imagine how beautiful the **rainbow** is.*

Richard Feynman



Similar may be said about *Quantum Interference*.

*Everybody knows that the interference does exist.
But it is not always easy to imagine
how it will work in a particular case.*

Yakov Azimov



- *Quantum physics* is probabilistic.
- *Classical physics* can be probabilistic as well (*Statistical physics*).
- **Essential difference:**
Classical Physics adds *probabilities*.
Quantum Physics adds *amplitudes*.
- **Important consequence:**
possibility of *interference* effects;
various *wave functions* may *mix*.
- **Impressive result:**
particles may *oscillate* in *time*,
transforming to *each other*.



Most Famous Examples:

- **Beauty oscillations** –
in decays of neutral **B**-mesons (@ *some μm*)
(coherent *oscillations* of several flavors are also possible)

Ya. Azimov, Phys Rev D **42**, 3705 (1990); Eur Phys J A **4**, 21 (1999)

- **Strangeness oscillations** –
in decays of neutral **K**-mesons (@ *$\sim 10\text{ cm}$*)

- **Neutrino oscillations** – 
(up to *hundreds km*, or even *astronomical* distances ! Solar ν)
Wide variety of *macroscopic* distances!

- Quite *macroscopic* manifestations of *quantum microscopic* effects !

- Same phenomenon may be *seen* in complementary variable – energy (mass in rest frame):
- It is seen here as *deformation* of BW peaks.



- Pure BW term: $|a (E - E_0 + i \Gamma/2)^{-1}|^2 = |a|^2 [(E - E_0)^2 + \Gamma^2/4]^{-1}$

- BW with background: $|B + a (E - E_0 + i \Gamma/2)^{-1}|^2$

$$= |B|^2$$

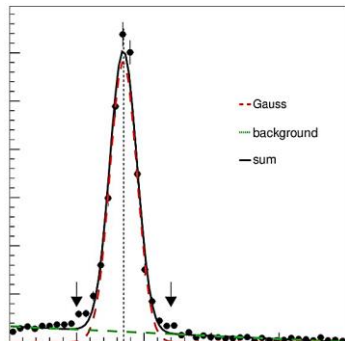
$$+ |a|^2 [(E - E_0)^2 + \Gamma^2/4]^{-1}$$

$$+ [2 |B a/ \cos\phi (E - E_0) + |B a/ \sin\phi \Gamma] \times [(E - E_0)^2 + \Gamma^2/4]^{-1}$$

may depend on E

interference term

role of interference depends on relative value & on relative phase ϕ of B & a ; it is linear in a , may change sign & be either positive or negative.



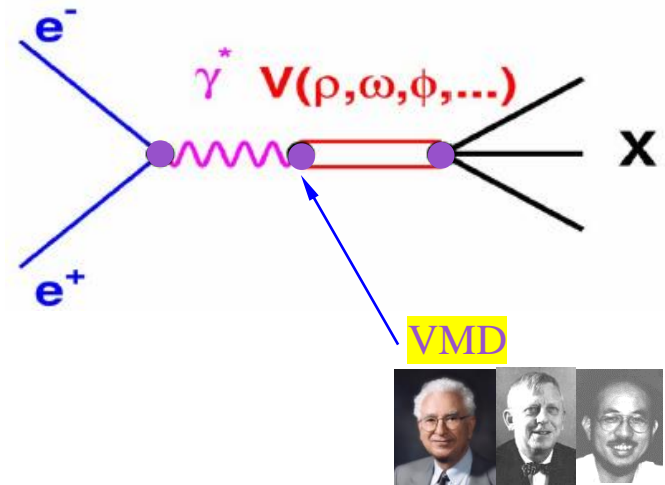
- @ small value of $|a/B|$ interference term may be more essential than proper BW contribution.
- Due to additional E -dependence, interference may change sign, provide either bump, or dip, or both.
- Bump &/or dip positions are, in general, shifted from true position of resonance.
- Same resonance may interfere differently in different decay modes.



- Rich source of examples, how **interference** works, is provided by reaction $e^+e^- \rightarrow \text{hadrons}$.

- Contributions with *same* final state are **coherent**; they all are produced through γ/Z & may **directly interfere**, if have *same* decay mode.

- Independent contribution of resonance is **BW**-peak, proportional to $\Gamma_{ee} \Gamma_X / \Gamma_{\text{tot}}$. **Interference** may **change** its form & intensity.



PDG $\Gamma(\rho) = 149.4 \text{ MeV}; \Gamma(\omega) = 8.5 \text{ MeV}; \Gamma(\phi) = 4.3 \text{ MeV}$

$\Gamma(\omega \rightarrow 3\pi) = 7.58 \text{ MeV}$

$\Gamma(\rho^0 \rightarrow 3\pi) = 0.015 \text{ MeV}$ **Isospin violated**

$\Gamma(\phi \rightarrow 3\pi) = 0.65 \text{ MeV}$ **Zweig rule violated**

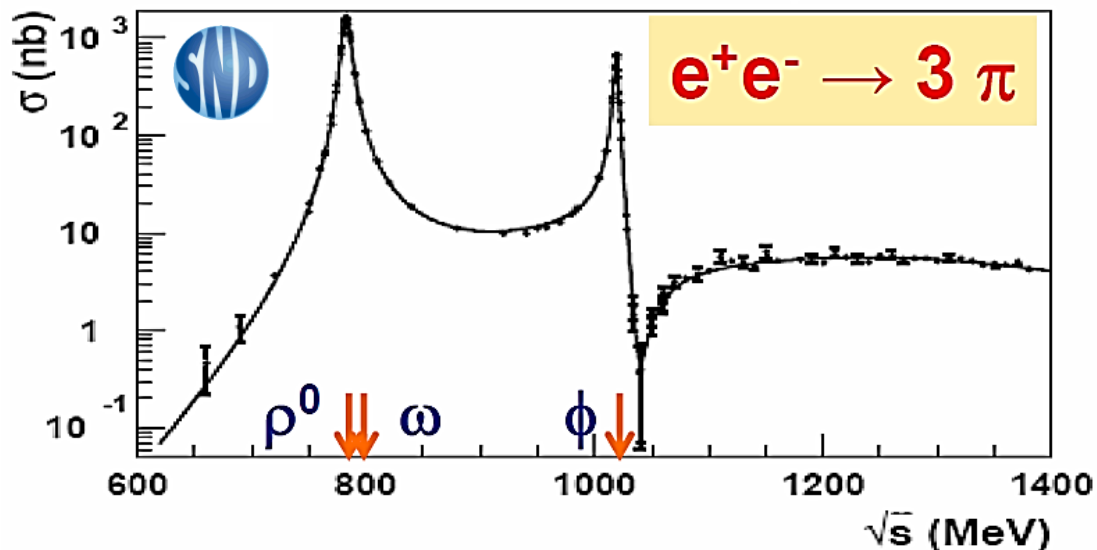


- Bkg near ϕ changes slowly



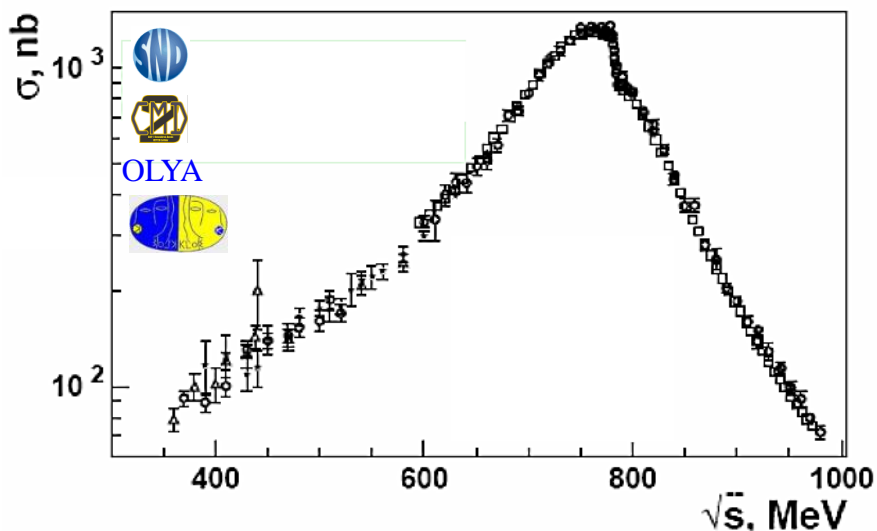
nearly standard **interference curve**,
instead of ϕ -peak:
both **bump** & **dip**,
each has form different from BW;
max/min different from ϕ -mass ρ .

- ρ -contribution here deforms ω -tails.
- Curve is fit with $\omega, \phi, \rho, \omega',$ & ω'' .



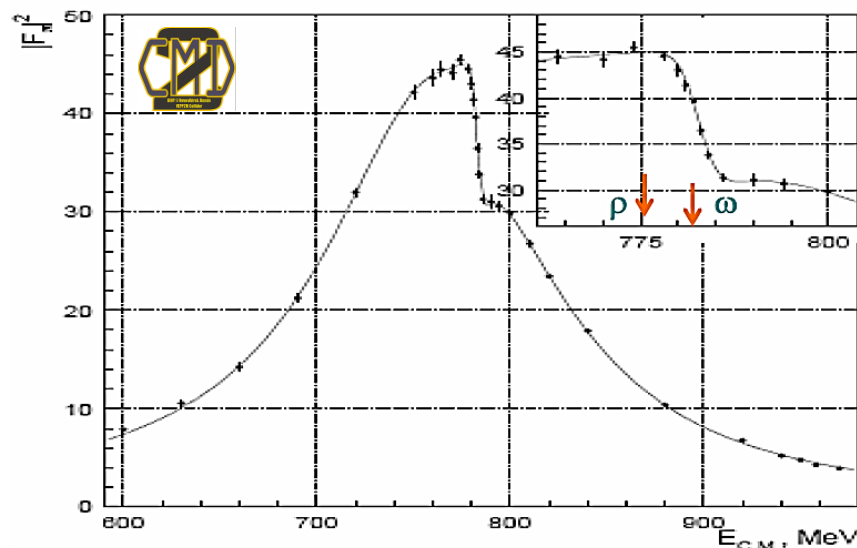
M. Achasov, Nucl Phys B Proc Suppl **162**, 114 (2006)

isospin violated
 $\Gamma(\rho^0 \rightarrow 2\pi) = 148 \text{ MeV}$
 $\Gamma(\omega \rightarrow 2\pi) = 0.13 \text{ MeV}$



• Most famous example of resonance interference.

- All above examples demonstrate **direct interference** of 2 resonances:
 all final particles can be decay products of any of the interfering resonances.
- Such kind of **interference** appears very efficient to search for rare decays of known resonances.
 may strongly deflect resonance manifestation from familiar BW peak.
- There can be other kinds of interference, where only some of final particles may come from any of 2 interfering resonances.



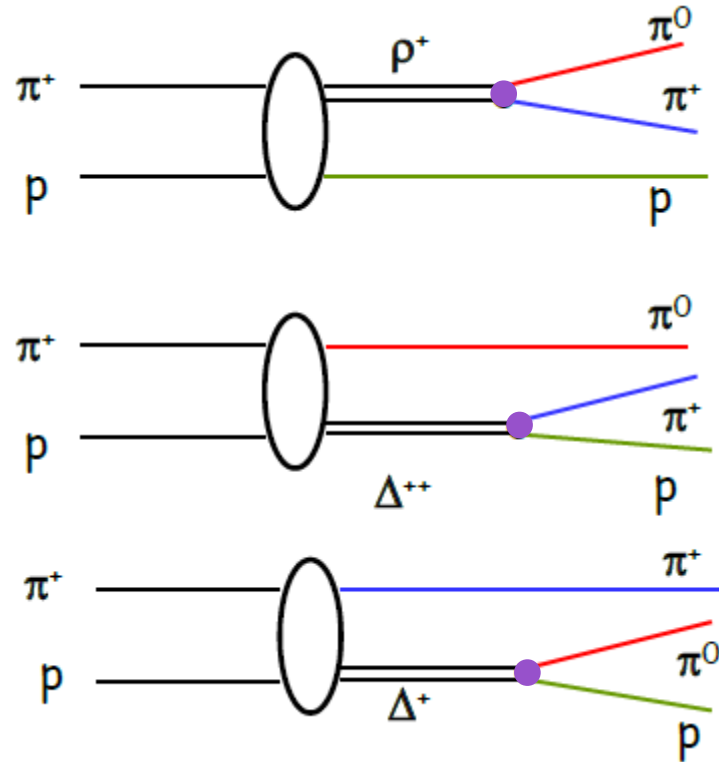
- Specifics of this case:
 rapidly decreasing bkg (ρ -peak);
 ρ - ω mixing may (& does) have complexity
 Y.I. Azimov, Eur Phys J A **16**, 209 (2003)
- ⇒ interference curve is strongly asymmetric:
 decrease, no increase.
- Opposite relative sign would reveal additional peak (case of $\eta\gamma$).
- Interference is only source of information on decay $\omega \rightarrow 2\pi$.



Rescattering Interference - I

- Different resonance configurations may produce same state of 3 or more particles.
- Such contributions are coherent & may interfere.

- Contributions depend on energies & momentum transfers; may shift & move positions of bumps/dips.

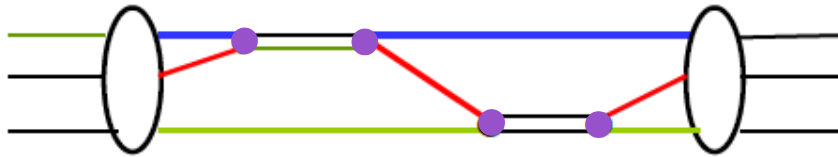


- Phenomenon is known since 60s.
- It was considered as *hindrance* to resonance studies.

- This interference was usually (& is still now) cut away.



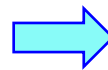
Rescattering Interference - II



- Name **rescattering** reflects similarity with **rescattering** in **3-particle** interactions: **one particle** changes its interaction partner.

- On other side, this kind of **interference** is like the famous case of **2 quantum slits**, since **one particle** refers **simultaneously** to **2 resonances**.
- Resonances are in **different systems** & may have **different quantum numbers** but final states, after resonance decays, should contain same particles.

• **2** resonances can **interfere**, only **if** final configurations are kinematically consistent.

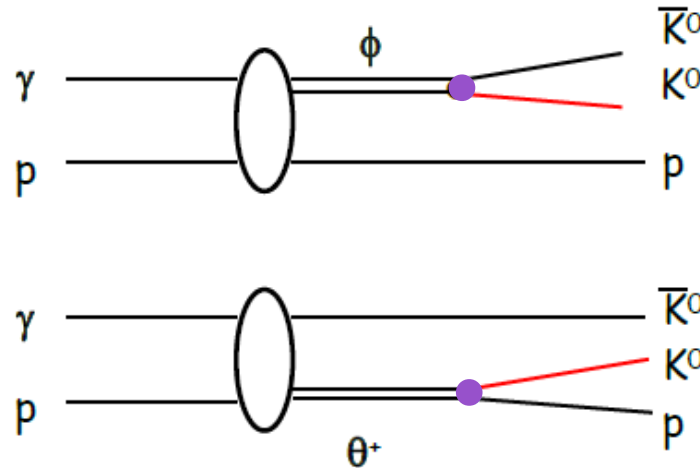


• This requires **limited intervals** of **total energy** & **momentum transfers**.

• Positions of **interference bumps/dips**, generally, depend on kinematic parameters & **move** with their changes (in difference with **true** resonance positions).

Rescattering Interference - III

- **Direct interference** of resonances has become efficient instrument actively used to study *rare decays* of **known** resonances.



- **Amaryan, Diakonov, Polyakov** have suggested to apply **rescattering interference** M. Amarian, D. Diakonov, & M.V. Polyakov, Phys Rev D **78**, 074003 (2008) for revealing small **new** resonance signals (amplification by *interference* with strong signal). CI=22

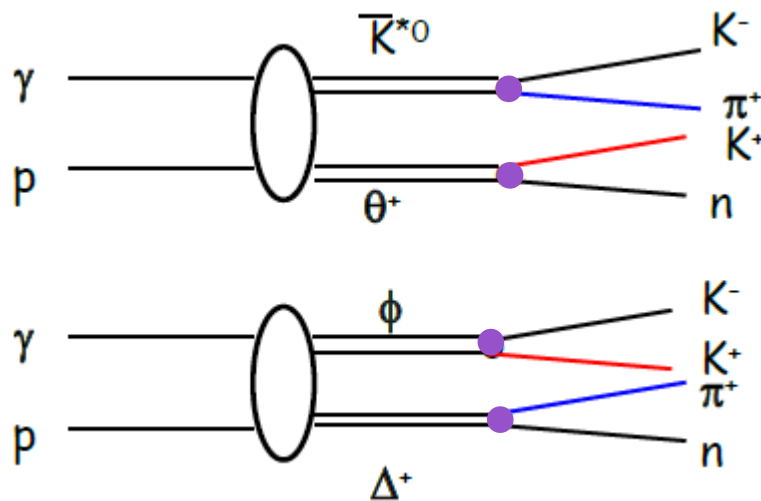


- Θ -signal may indeed be **small**, if its production is new kind of *hard processes*. Y.I. Azimov, K. Goeke, & IIS, Phys Rev D **76**, 074013 (2007) CI=17



Rescattering Interference - IV

- *Final states* with **> 3** particles admit more complicated cases of *rescattering interference*.



- Example of **4-particle rescattering interference** that may enhance small θ^+ -contribution (suggested by Amaryan).

Intermediate Summary

- *Interference* of resonances (in energy representation) has same origin as known particle oscillations (in space-time representation).
- Small resonance contribution may be *amplified* & revealed due to its *interference* with high background (e.g., another resonance).
- Manifestation of *interfering* resonance may be very different: **bump**, or **dip**, or **both**; may depend on decay mode. Positions of **bump/dip** are, in general, *shifted* from *true* position of resonance.
- Form of “resonance” curve essentially depends on properties of background & on Res-Bkg *relative phase*: it may be *symmetric*, or *antisymmetric*, or *strongly asymmetric*.

- *Direct interference* is actively used now as important instrument for resonance studies. Some rare decays of well - established resonances are known **only** due to *interference* manifestations.
- *Rescattering interference* of resonances may be very useful as well: to amplify small resonance signals, especially with **new quantum numbers**; to study **production mechanisms** of known resonances.
- *Interference of resonances* looks to be worth of more detailed studies, both experimental & theoretical.



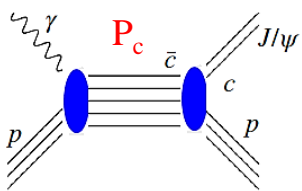
Back to Mainstream

Alternative Solution for GlueX J/ψ Data



Recipe for Possible Interpretation of Dip

IIS, W.J. Briscoe, E. Chudakov, I. Larin, L. Pentchev, A. Schmidt, R.L. Workman,
 Phys Rev C **108**, 015202 (2023) CI=2



- Experimental total Xsec of inelastic binary reaction: $\sigma_t = \int_0^{2\pi} \int_0^\pi \frac{d\sigma}{d\Omega} \sin\theta \, d\theta \, d\phi$

θ & ϕ are J/ψ polar & azimuthal production angles.

- Phenomenological total Xsec: $\sigma_t = \frac{\pi}{4k^2} \sum_{J=0}^{\infty} (2J+1) |f|^2$

k is momentum of photon in CM
 J is total angular momentum
 $(2J+1) = 1$ for S -wave
 $= 3$ for P -wave

Using Landau-Livshitz normalization

- Partial Amplitude: $f = b + R \cdot \exp(2i\alpha)$

α is relative phase shift

It comes from fit of total Xsec

There is 1 free parameter for interference α

IIS, A.V. Kravtsov, & M.G. Ryskin, Sov J Nucl Phys **40**, 274 (1984)

- Non-Res: $b = \sqrt{A} q + \sqrt{B} q^3$

q is momentum of vector meson in CM

There are 2 free parameters for background A & B

IIS, L. Pentchev, & A.I. Titov, Phys Rev C **101**, 045201 (2020)

- Relativistic BW: $R = \frac{2\Gamma M}{[(M)^2 - s] - i\Gamma M} X$

M & Γ are mass & energy independent width
 $(P_c$ is too narrow)

- Partial Width: $X = \frac{\sqrt{\Gamma(\gamma+p) \Gamma(J/\psi+p)}}{\Gamma} = \sqrt{X(\gamma+p) X(J/\psi+p)}$

There are 3 free parameters for resonance M , Γ , & X

$\Gamma(\gamma p)$ & $\Gamma(J/\psi p)$
 are partial decay widths of
 $P_c \rightarrow \gamma p$ & $P_c \rightarrow J/\psi p$.



How Interference Works

Ya. Azimov, J Phys G 37, 023001 (2010)



CI=18

- *Interference* of resonances has same *quantum nature* as oscillations of particles, though they are observed in complementary variables – energy (mass) for former, or time for latter.

- *Interference* contribution is linear in both b & R .

Its relative role depends on R/b .

@ *small* R/b , *interference* may appear more important than BW contribution itself.

This can be considered as *amplification* of small resonance signal by *interference* with large *background*.

$$f = b + R \cdot \exp(2i\alpha)$$

- *Interference* contribution depends on *relative phase* α between b & R .

- *Interference* may be either *positive (constructive)* or *negative (destructive)*.

- In comparison with BW contribution, *interference* may have additional *energy dependence* & may decrease with energy slower than proper BW contribution.

- *Background* itself, b & α , also may depend on energy.

As a result, *background* may appear different in regions of *constructive* & *destructive interference*, & relative role of these regions may be very different, up to full vanishing of one of them.

Thus, presence of *interference* may provide either *bump*, or *dip*, or *both*.

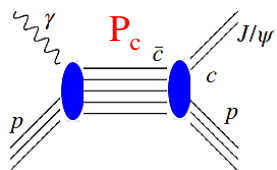
Positions of *bump* &/or *dip* are, generally, *shifted* from true position of *resonance*.

- Resonances can *interfere* differently in different decay channels, @ least due to different properties of corresponding *backgrounds*.



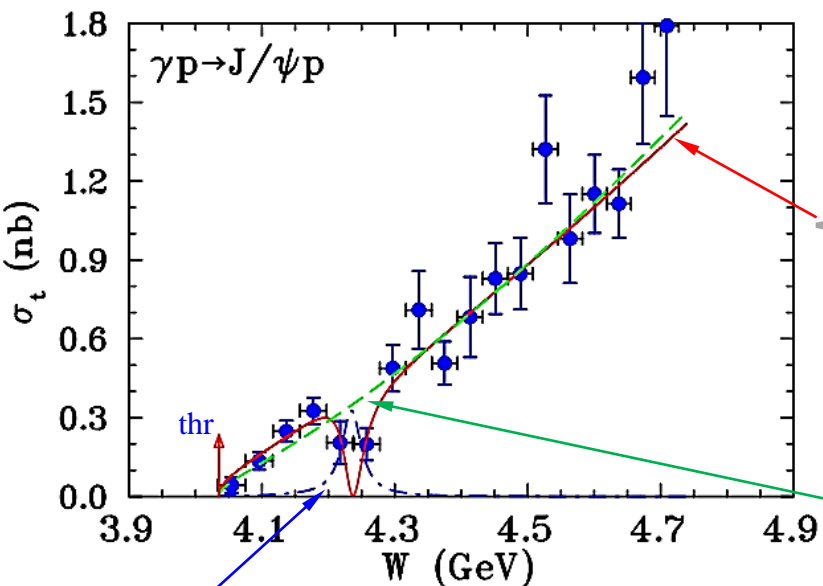
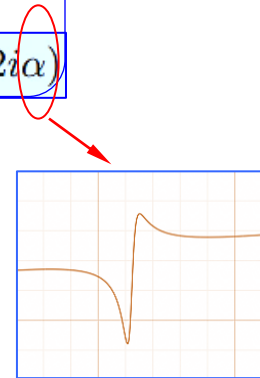
Alternative Solution for Data

IIS, W.J. Briscoe, E. Chudakov, I. Larin, L. Pentchev, A. Schmidt, R.L. Workman, Phys Rev C **108**, 015202 (2023)



- We suggested to apply *rearrangement interference* for revealing *faint* resonance signals (*amplification* by *interference* with *strong* background signal).
- Relative phase α leads to *constructive* (*bump*) or *destructive* (*dip*) *interference* for particular **PW**.

$$f = b + R \cdot \exp(2i\alpha)$$



Resonance: $\chi^2/ndf=11.99/12=1.00$

$M = 4235 \pm 8$ MeV

$\Gamma = 35.4 \pm 8.2$ MeV

$X = 0.023 \pm 0.005$

$\alpha = 40.8 \pm 5.7$ deg

Resolution ~ 6 MeV

Background:

$A = 0.00251 \pm 0.00046$ nb GeV/c


$B = 0.00688 \pm 0.00083$ nb/GeV/c

No Resonance: $\chi^2/ndf=19.74/16=1.23$

$A = 0.00183 \pm 0.00040$ nb GeV/c

$B = 0.00766 \pm 0.00077$ nb/GeV/c

- *Dip* position does not correspond to *real mass* of $P_c(4312)^+$.
- It may depend on reaction *mechanism* (including *cusps* (*open charm*)) & background choices.

- If “*bump*” is imposed on  data “*by hand*” (consider **7th - 9th** energy values up from threshold), qualitative description of data up to $W = 4.35$ GeV is possible, but with higher χ^2 , if our fit form is used.

- Obtained mass in our analysis is almost **77 MeV** below  determination, but it cannot exclude that this is $P_c(4312)^+$.

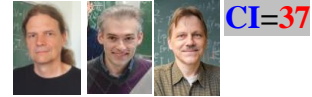


Cusp Effect



Deciphering Mechanism of Near-Threshold J/ψ Photoproduction

Meng-Lin Du, V. Baru, Feng-Kun Guo, Ch. Hanhart, U.-G. Meissner, A. Nefediev, & IIS, Phys Rev C **106**, 015202 (2022)

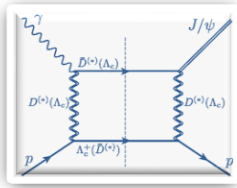
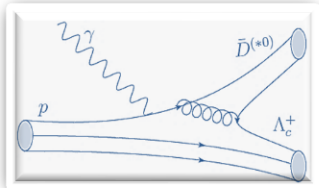


CI=37

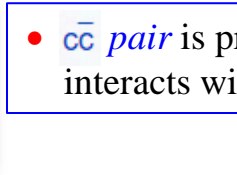
- It was shown that *fluctuation of photon into open charm* $\gamma p \rightarrow \Lambda_c \bar{D}$ is preferable than into *Charmonium J/ψ* . K. Boreskov, A. Capella, A. Kaidalov, & J. Tran Than Van, Phys Rev D **47**, 919 (1993)



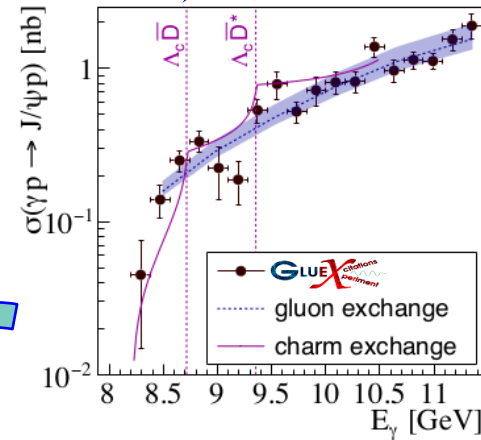
Cusp effect is visible & in agreement with ~~GLUEX~~ ^{collaboration experiment}



$c\bar{c}$ pair is produced by $1g$ & interacts with proton.

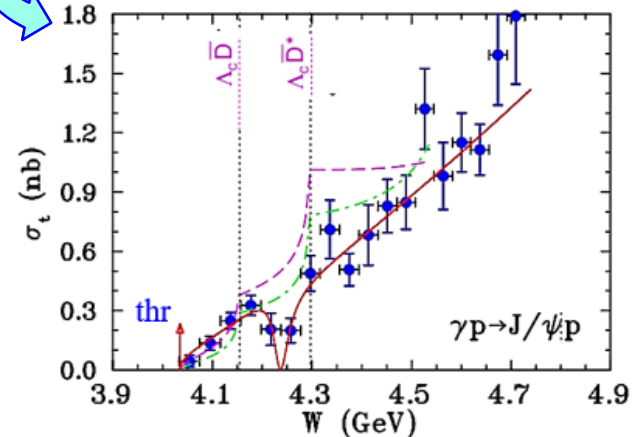


$c\bar{c}$ pair is produced by *photon* via VMD & interacts with *proton* through $2g$ exchange.



- These *two mechanisms* act simultaneously. Assuming there is only *first one*, then key consequence: *threshold cusps* !
- There is no fit to ~~GLUEX~~ ^{collaboration experiment} data.

- One should study *two-component* problem accounting for *interference* between these *two components*.
- Effect of *charm* exchange is smaller than *gluon* exchange.
- Glueon* contribution can be strongly *suppressed* due to “*young*” effect.



E.L. Feinberg, Sov Phys Usp, **23**, 629 (1980)
Courtesy of Misha Ryskin, July 2020



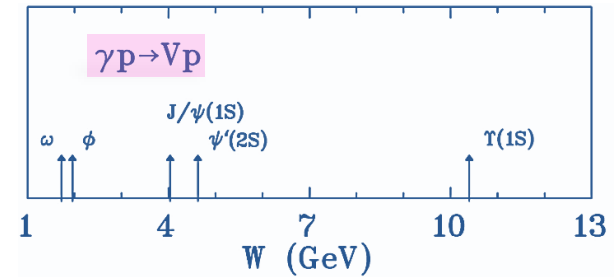
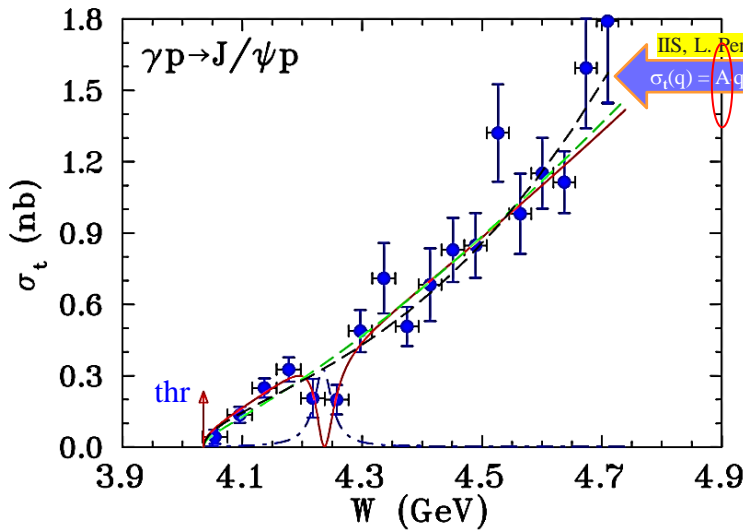
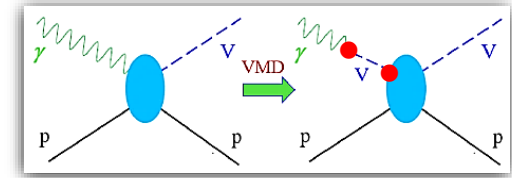
Vector Meson – Nucleon Scattering Length

Joint HEPD-THD Seminars (Thursdays) April 6, 2023

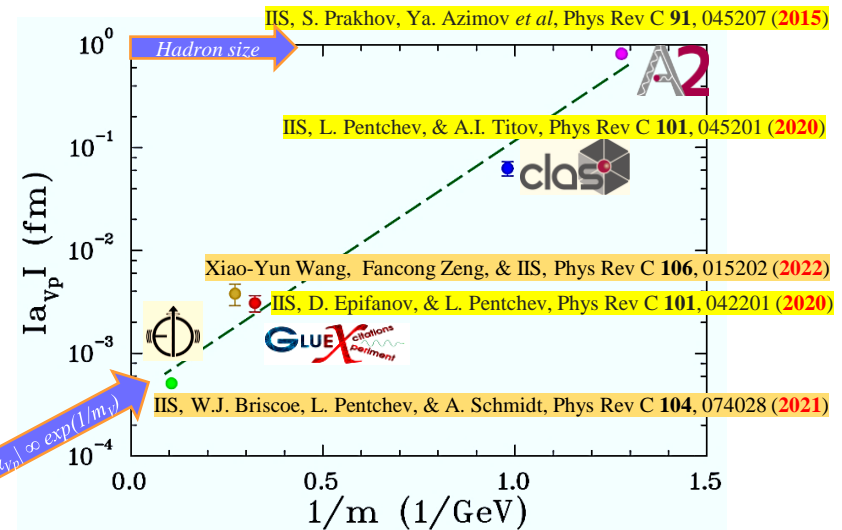


Vector Meson – Nucleon SL

- Due to *small size* of “young” V vs “old” V , measured & predicted SL is very small.
- V created by photon @ threshold then most probably V is not formed completely & its radius is smaller than that for normal (“old”) V .
- Therefore, one observes stronger suppression for Vp interaction.



	A	$ a_{J/\psi p} $
	[nb/MeV/c]	[mfm]
2019	0.46 ± 0.16	$3.08 \pm 0.55 \pm 0.42$
2023	0.43 ± 0.10	2.88 ± 0.34



- $p \rightarrow V$ coupling $\bar{q}q$ is proportional to a_s & separation of corresponding quarks.
- This separation (in zero approximation) is proportional to $1/m_V$.

JPAC Contribution



- Originally PWA arose as technology to determine amplitude of reaction via fitting scattering data.

⇒ That is *non-trivial mathematical problem* – looking for solution of **ill-posed** problem following to Hadamard & Tikhonov.
[number of equations less than number of unknown quantities]

⇒ There are two main technologies to look for solution:

(i) *least-squares minimization* of functions which are linear in unknown parameters, χ^2 &

(ii) *likelihood measures goodness* of fit of statistical model.

[*Minimizing χ^2 is equivalent to maximizing (log) likelihood*
just case *not small statistics*]

⇒ Model *independent* treatment or data *driven* treatment.



Roger Cotes



Sir Ronald Aylmer Fisher

- Resonances appeared as by-product

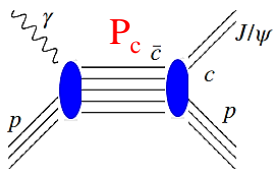
[bound states objects with definite quantum numbers, mass, lifetime, & so on].

- Standard PWA

⇒ Reveals only wide Resonances, but not too wide ($\Gamma < 500$ MeV)
& possessing not too small BR (BR > 4%).

⇒ Tends (by construction) to miss narrow Resonances with $\Gamma < 20$ MeV.



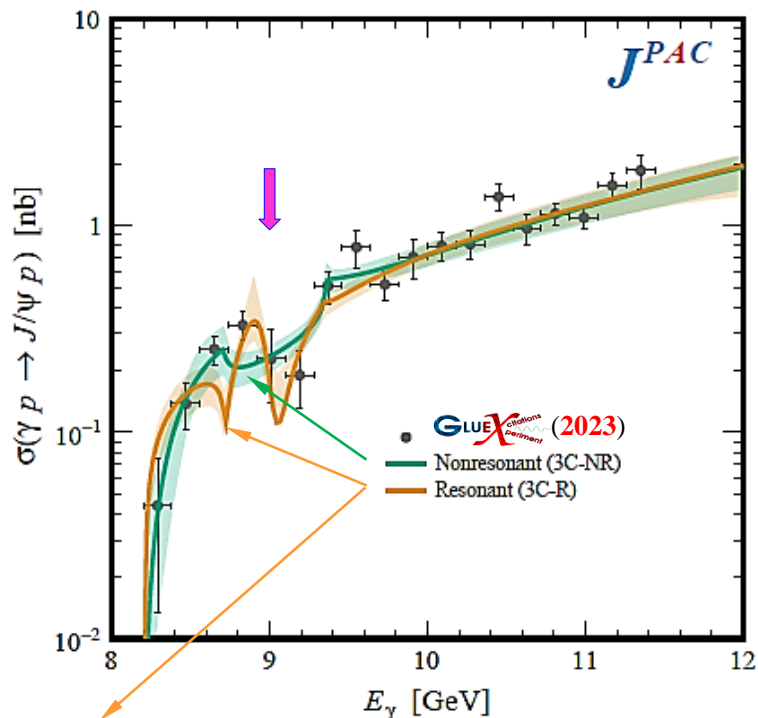


JPAC for J/ψ Photoproduction

D. Winney *et al*, Phys Rev C **108**, 054018 (2023)



CI=4



$|a_{\psi p}| \lesssim 0.2$ fm at a 90% CL

Much-much larger than we found.

Pole position

(no uncertainties)

$M = 4211$ MeV

$\Gamma = 48$ MeV

101 MeV below LHCb

24 MeV below us

Phenomenological model based on s-channel PW expansion ($1 < 3$):

- Global fit of both J/ψ & D $d\sigma/dt(t)$ & Hall D $\sigma_{tot}(E_\gamma)$
 - (1C) J/ψ interaction
 - (2C) J/ψ p interaction & $\bar{D}^* \Lambda_C$
 - (3C-NR) J/ψ p, $\bar{D} \Lambda_C$, $\bar{D}^* \Lambda_C$ (non-res soln)
 - (3C-NR) J/ψ p, $\bar{D} \Lambda_C$, $\bar{D}^* \Lambda_C$ (res soln)

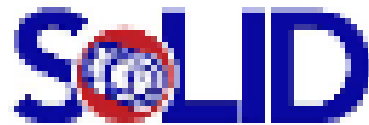
No stat. significant preference:

- $E_\gamma = 9$ GeV structure requires sizable contribution from *open charm*.
- Severe violation of **VMD** & factorization *not excluded*.
- s-channel resonance *not excluded*.
- t-enhancement maybe due to *proximity* to threshold (s-wave only).

Precise measurements critically important to disentangle *reaction mechanism* @ threshold.

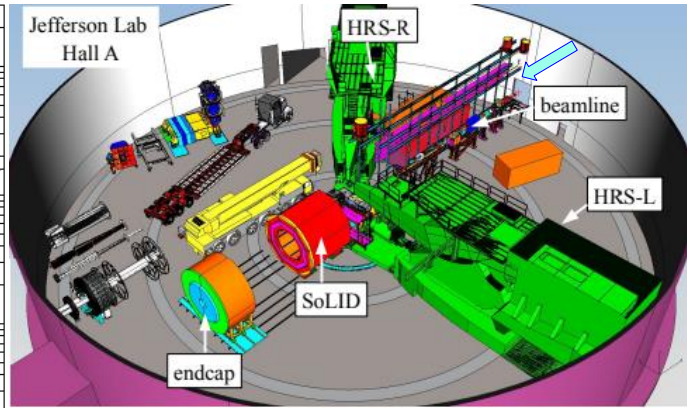
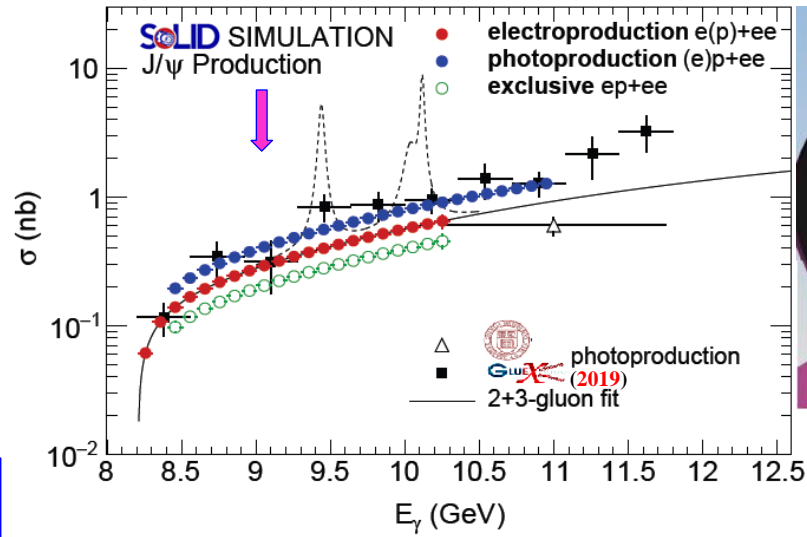


SOLID Contribution

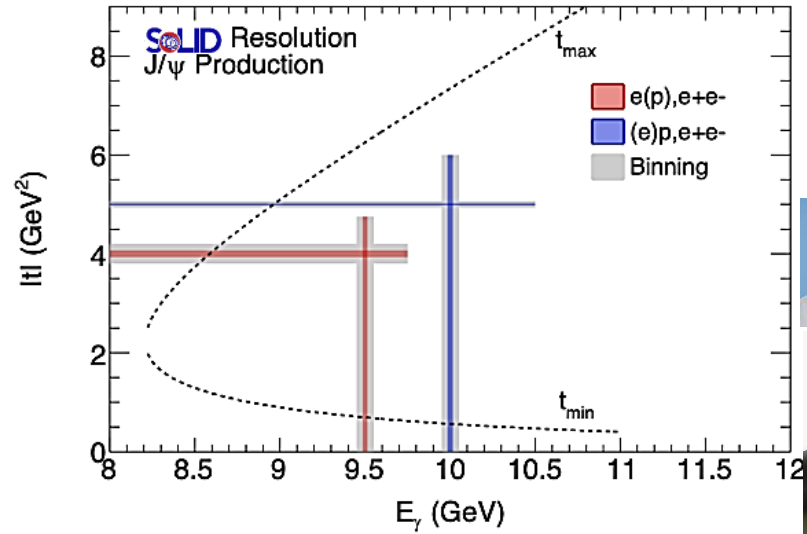


Jefferson Lab Hall A for more J/ψ Photoproduction

Bump Hunting



• Nominal luminosity for SoLID- J/ψ is 50 days @ $10^{37} \text{ cm}^{-2} \text{ s}^{-1}$



• Only downside is we, probably, must wait 10 years for such results.

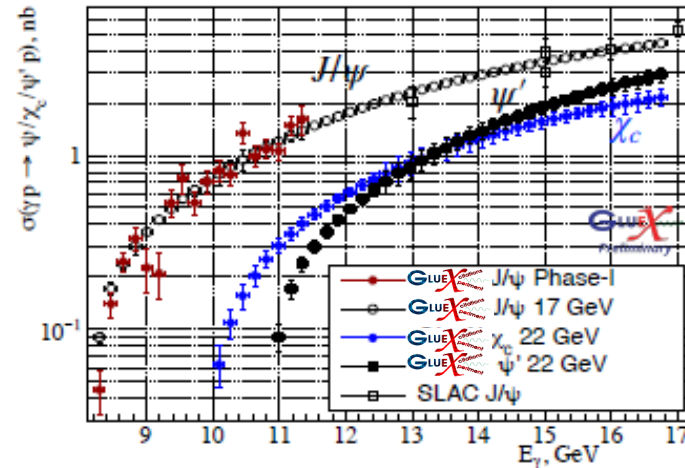
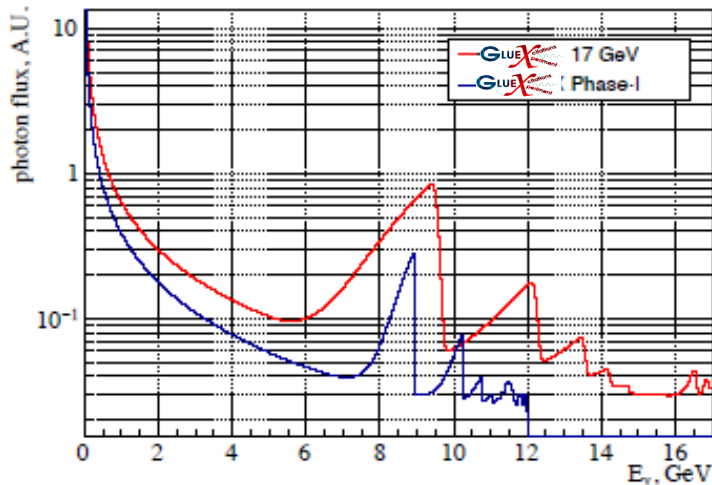
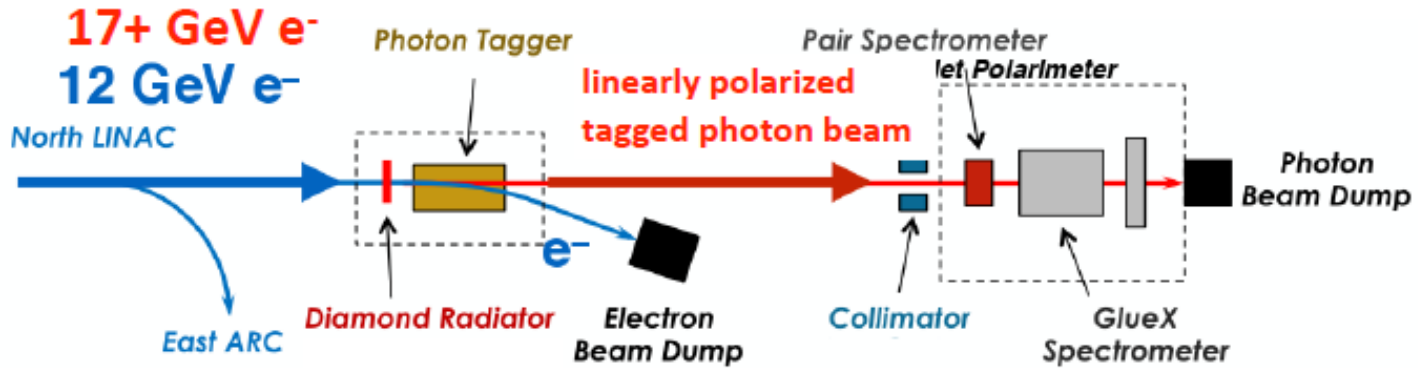
J. Arrington et al, arXiv:2209.13357 [nucl-ex]



JLab at High Energies



Hall D Apparatus with 17+ GeV electron beam



- Moving end point from 12 GeV to 17+ GeV:
 - higher flux (and polarization) toward higher energies, while low energies less affected (no load on detectors)

Courtesy of Lubomir Pentchev, 2023



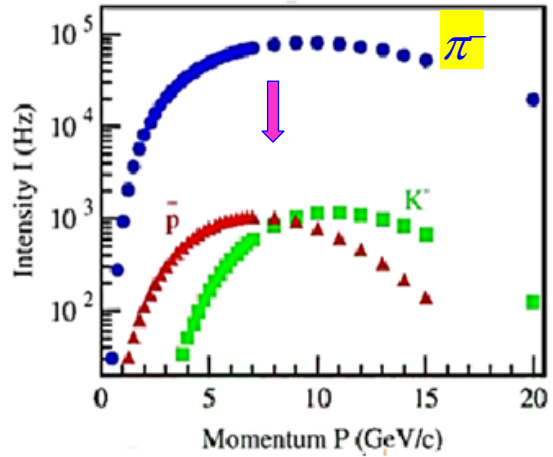
J-PARC Contribution



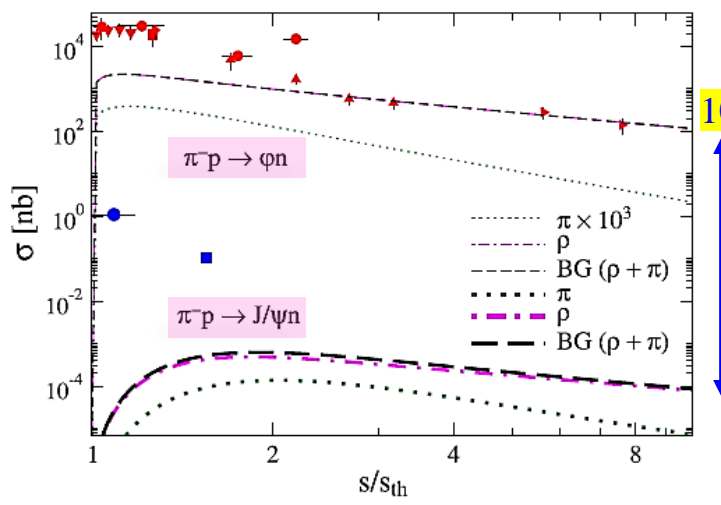


Bump Hunting

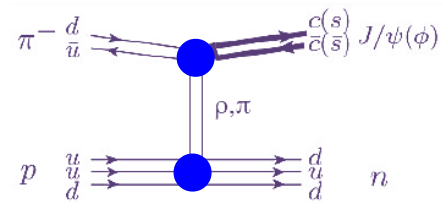
Hidden Charm Production by High-p Experiment





- High-p plans to detect J/ψ to e^+e^- & $\mu^+\mu^-$ pairs using E50 spectrometer.
- High-p can use incident $P = 2 - 20$ GeV/c beam from $\pi 20$ beamline.
- One can measure J/ψ production @ $P = 8 - 10$ GeV/c.
- $W_{thr} = 4$ GeV ($P_{thr} = 8.06$ GeV/c).
- Momentum bite is expected to be $\pm 3\%$.



$10^6 g_{\phi\rho\pi} > g_{J/\psi\rho\pi}$



- New High-p measurement allows to understand dynamics of $c\bar{c}$ production @ threshold.
- It is free from VMD & allows to determine $J/\psi p$ SL independently on .
- It allows to look for effect of .

S.H. Kim, H.C. Kim, & A. Hosaka, Phys Lett 763, 358 (2016)











Summary





SUMMARY

- Interpretations of published  &  cases considered *constructive interference* between resonance & non-resonant background for particular **PW**,
In this work, we relax this requirement & consider a possibility of *destructive interference* which *nobody did before*.
- Here we have shown that *resonance-like* structure is “*plausible*” in new  data, in energy region close to low-mass  pentaquark.
- Statistics of new  data is not sufficient to draw definite conclusion.
- Predictions of *open charm* model reasonably agree with new  data.
Interference between *open charm* & *gluon exchange* may (by some accident) produce *dip* but there is room for *resonance*.

While not evident in  data, one cannot exclude that we have all 4  P_c resonances together with *open charm* & *gluon exchange* (*gluon* contribution can be strongly suppressed due to “*young*” effect).

- It seems that analysis of full  -I &  -II &/or  may help.
It will be good to shrink significantly *uncertainties* (to increase statistics by a factor of **10+**) & make much larger number of cross sections (by factor of **2+**).
Polarized measurements are important contribution for model independent PWA.



may help to solve this puzzle.

Thanks





10/8/2023

Joint Seminar HEPD-THD, PNPI, Gatchina, Russia, October 2023

Igor Strakovsky 44

